

# Charge-transfer features in the exact Kohn-Sham potential of ensemble density functional theory



MAX-PLANCK-GESELLSCHAFT



האוניברסיטה העברית בירושלים  
THE HEBREW UNIVERSITY OF JERUSALEM



**Matt Hodgson, Eli Kraisler, Axel Schild and Hardy Gross**

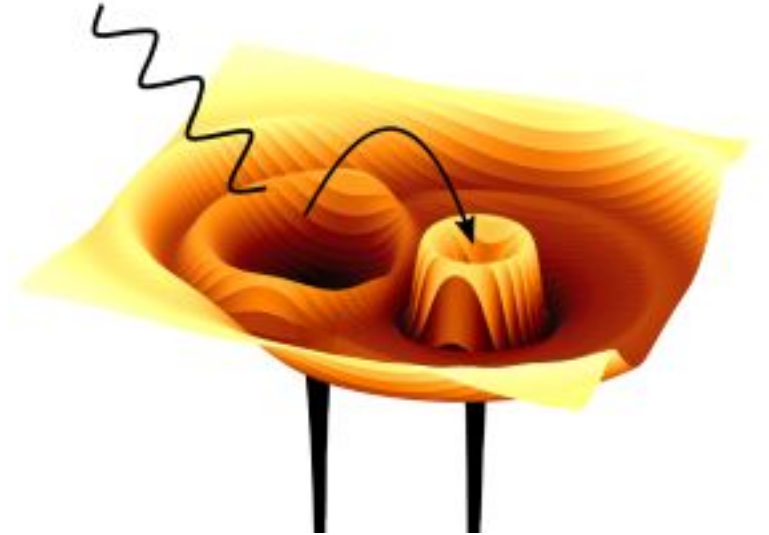
Max Planck Institute of Microstructure Physics, Halle, Germany

The Hebrew University of Jerusalem, Jerusalem, Israel

ETH Zürich, Zürich, Switzerland

# Charge transfer in a molecule

When the electrons within a molecule are excited charge is transferred among the atoms



# So what's the problem?

The more electrons you have, the more interactions...

...the exact, full-correlated many-body wavefunction quickly becomes impossible to calculate



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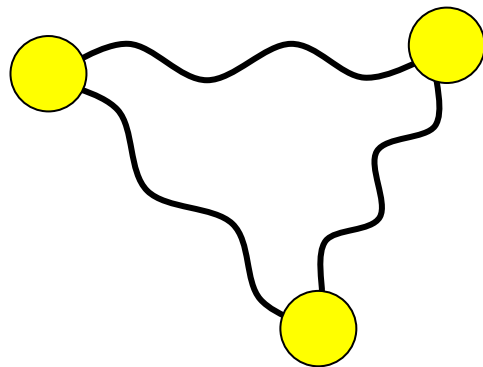
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$$\Psi(x_1, x_2, x_3)$$

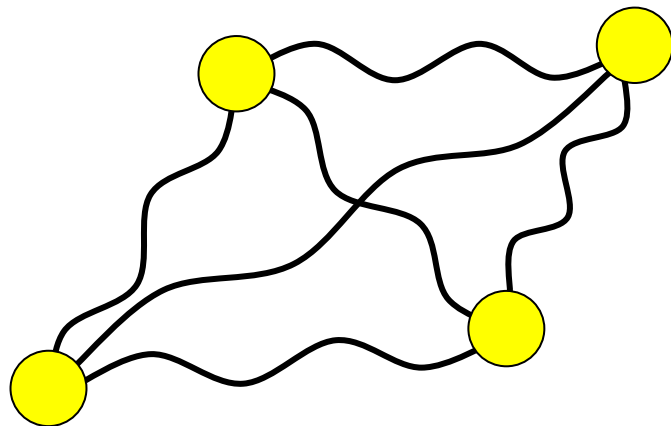


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$$\Psi(x_1, x_2, x_3, x_4)$$

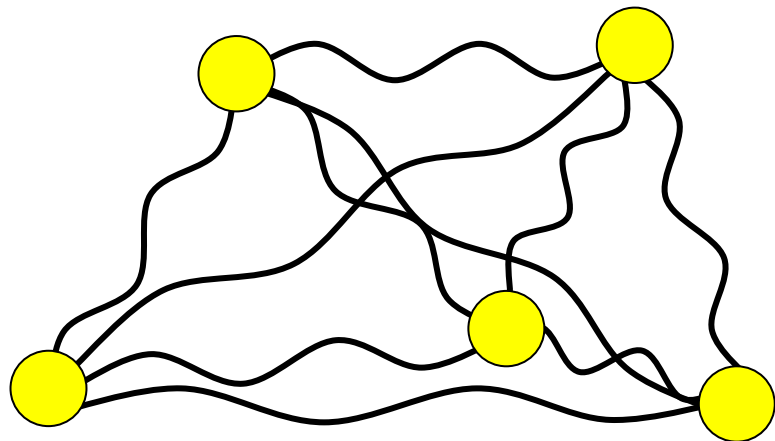


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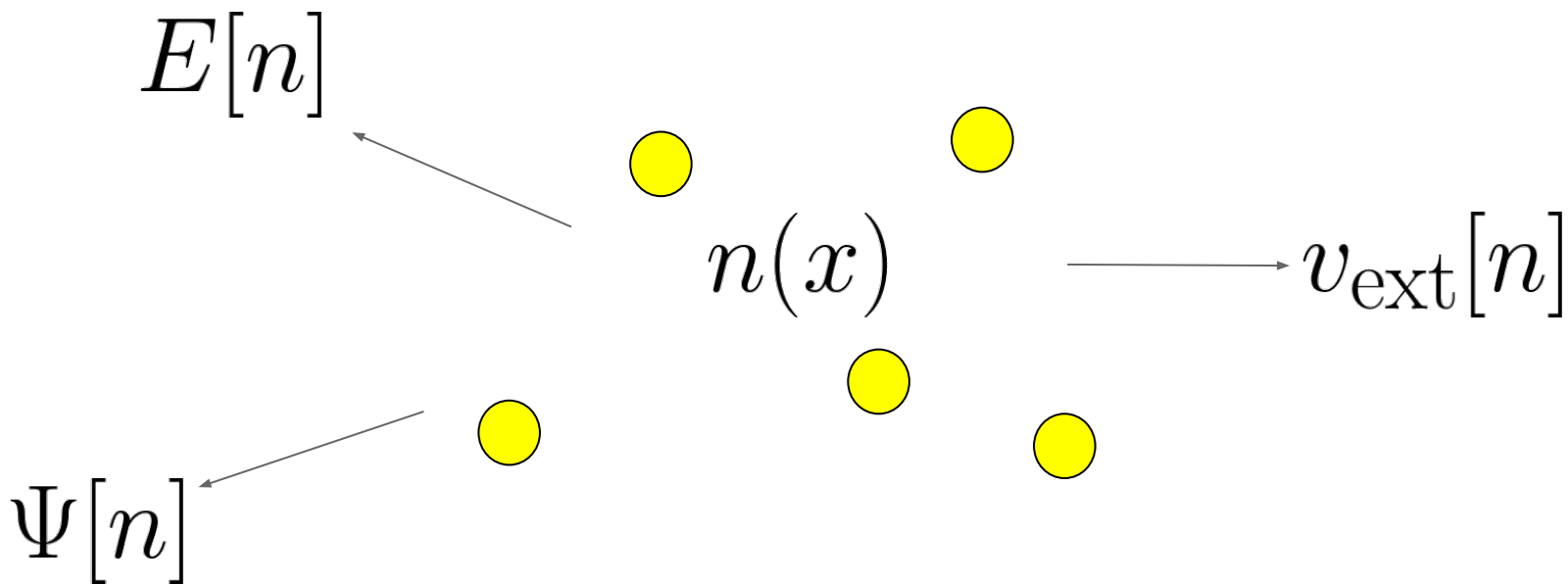
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$$\Psi(x_1, x_2, x_3, x_4, \dots)$$



# Density functional theory

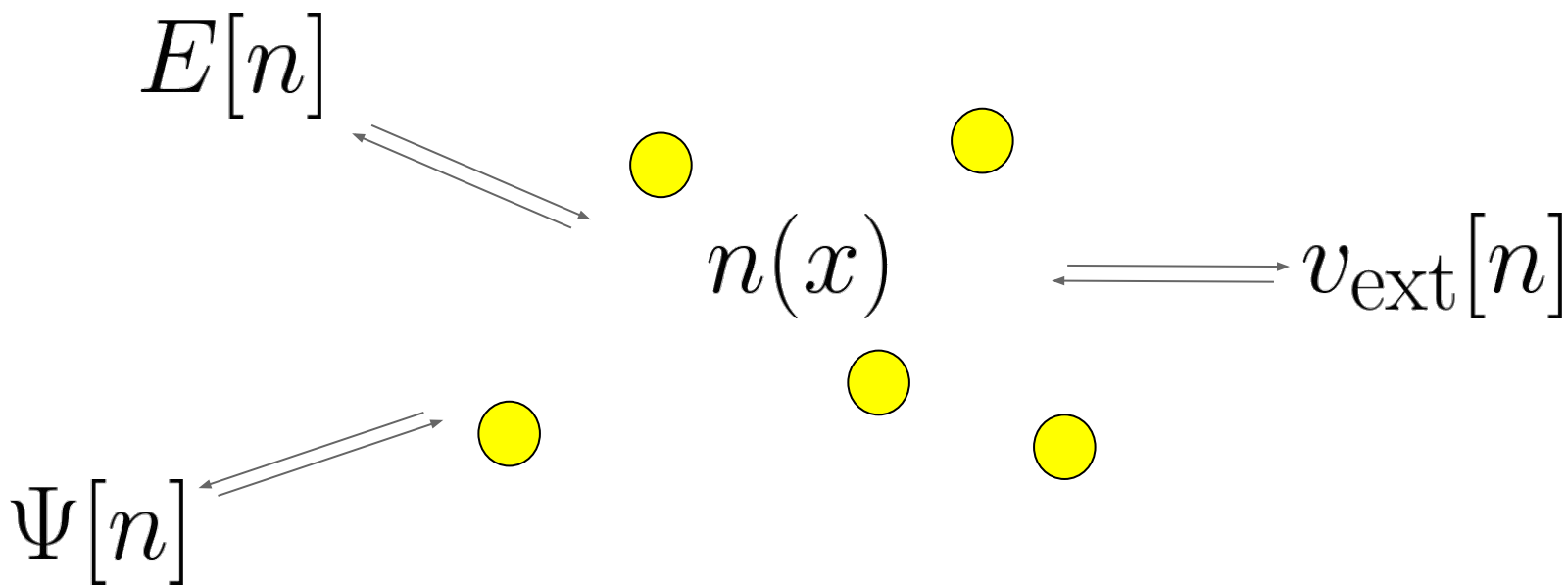
Replace the wavefunction with the *density* – P. Hohenberg and W. Kohn Phys. Rev. **136**, B864





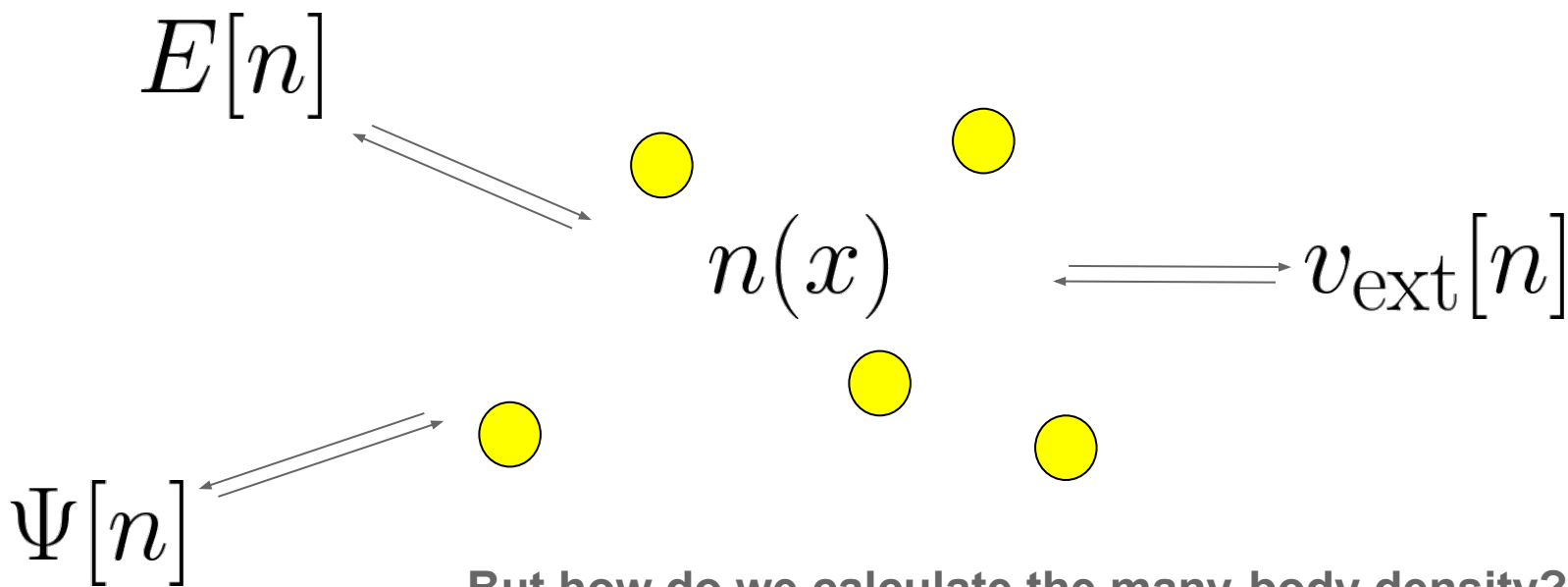
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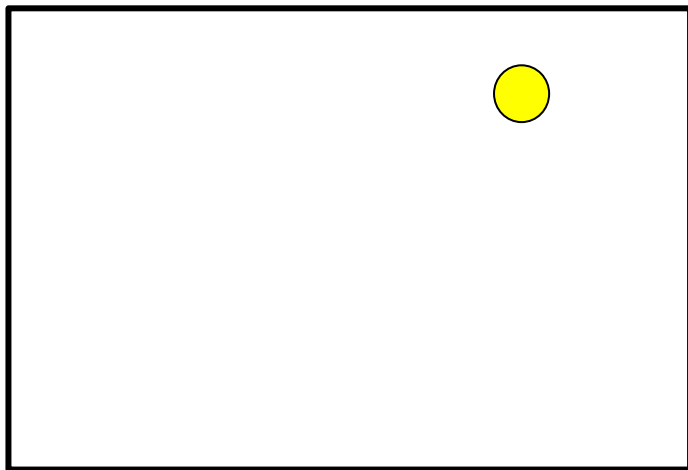
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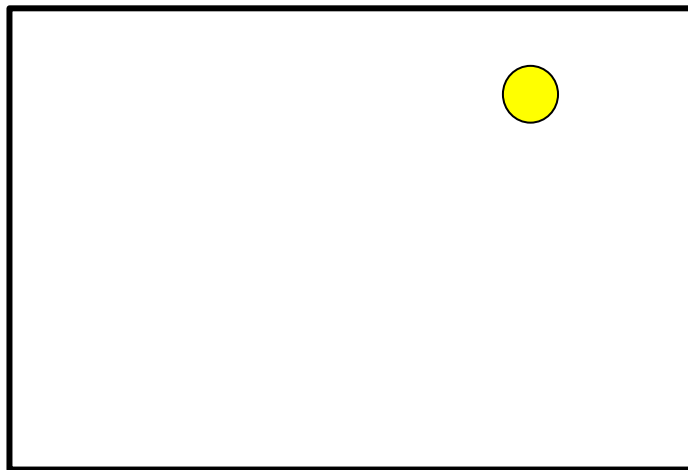


But how do we calculate the many-body density?

## Many-body system



## Kohn-Sham system

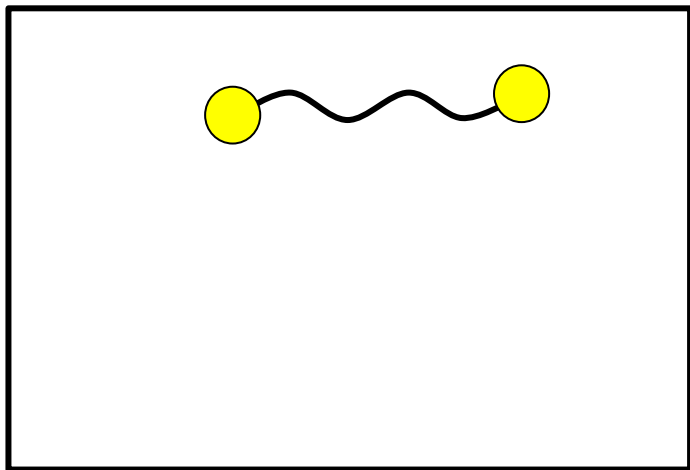


**The electron densities are the same**

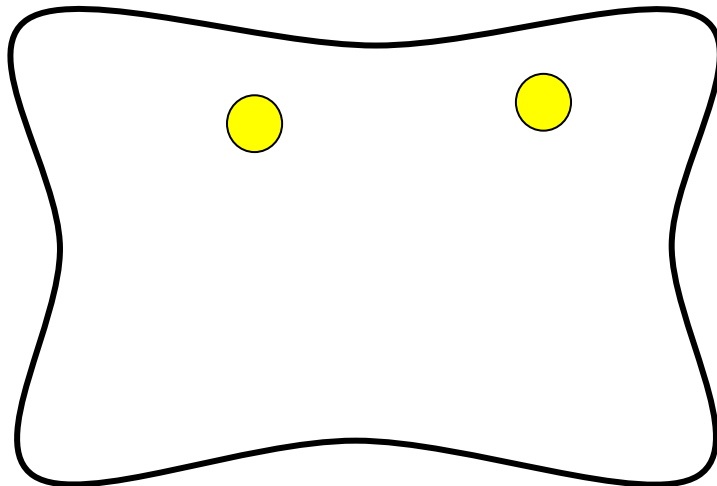
Computational scaling:  $k^1$

Computational scaling:  $1 \cdot k$

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## Kohn-Sham system

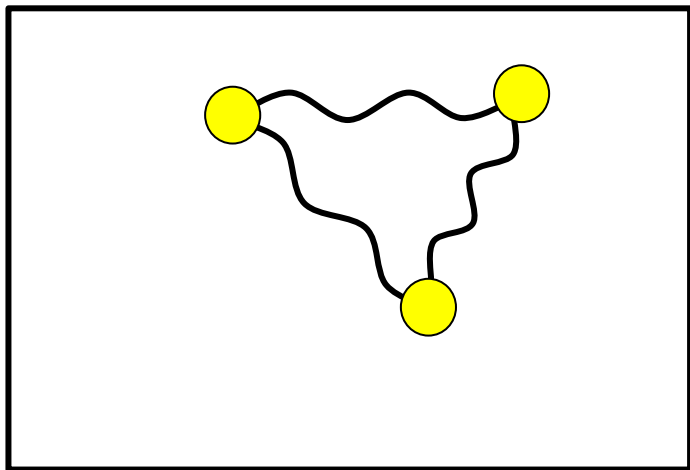


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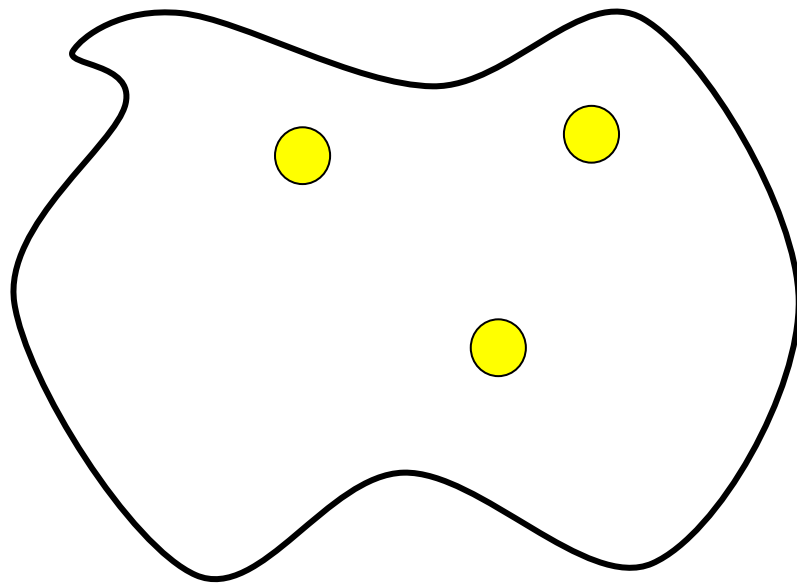
Computational scaling:  $k^2$

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## Many-body system



## Kohn-Sham system

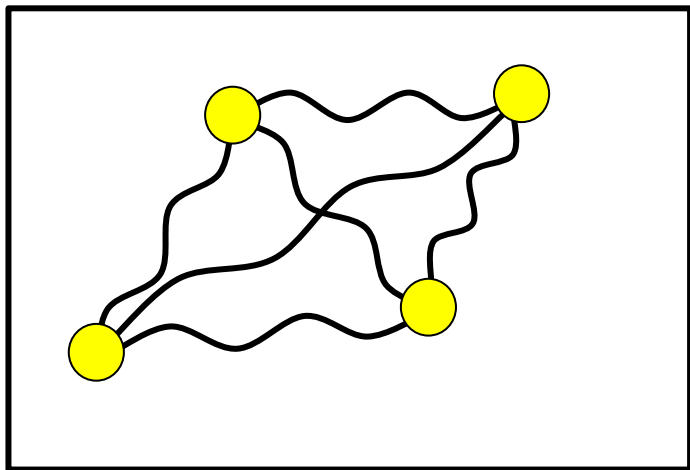


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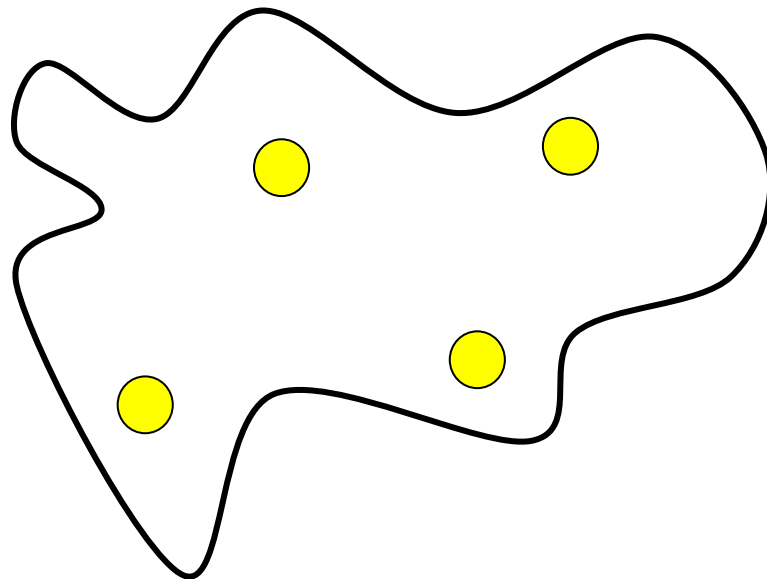
Computational scaling:  $k^3$

Computational scaling:  $3 \cdot k$

## Many-body system



## Kohn-Sham system

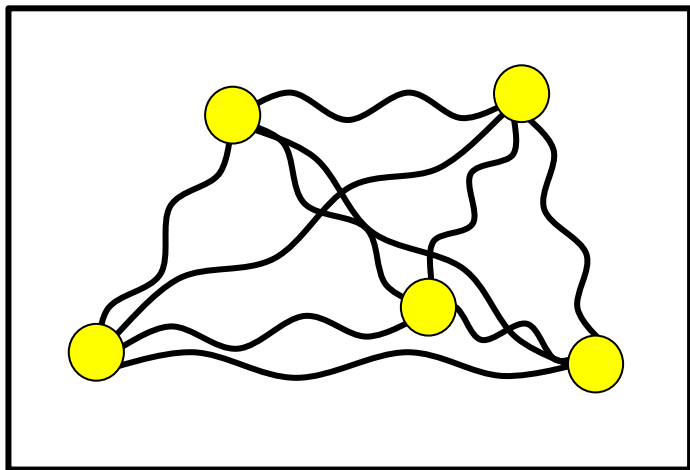


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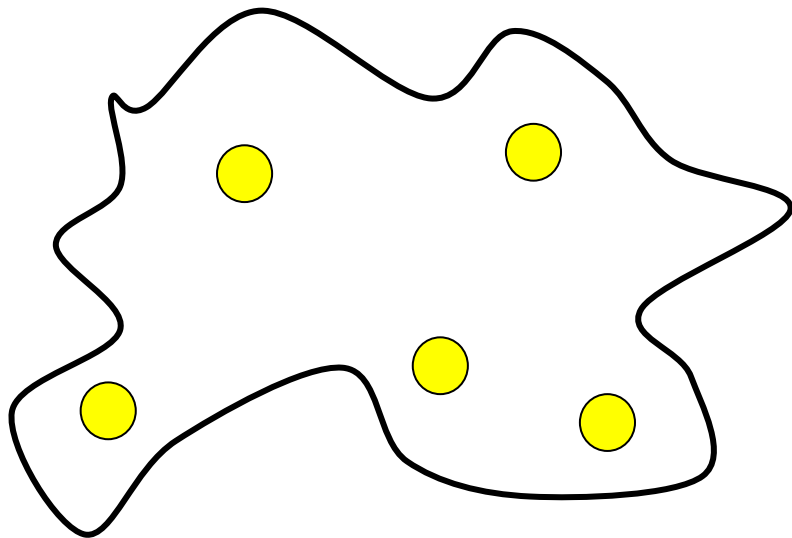
Computational scaling:  $k^4$

Computational scaling:  $4 \cdot k$

## Many-body system



## Kohn-Sham system



**The electron densities are the same**

Computational scaling:  $k^5$

Computational scaling:  $5 \cdot k$

# Kohn-Sham theory

Auxiliary system of fully non-interacting electrons which experience a local effective potential – the Kohn-Sham potential

$$\left( -\frac{1}{2} \frac{d^2}{dx^2} + v_s(x) \right) \phi_i(x) = \varepsilon_i \phi_i(x) \quad n(x) = \sum_i |\phi_i(x)|^2$$

$$v_s(x) = v_{\text{ext}}(x) + v_{\text{H}}(x) + v_{\text{xc}}(x)$$



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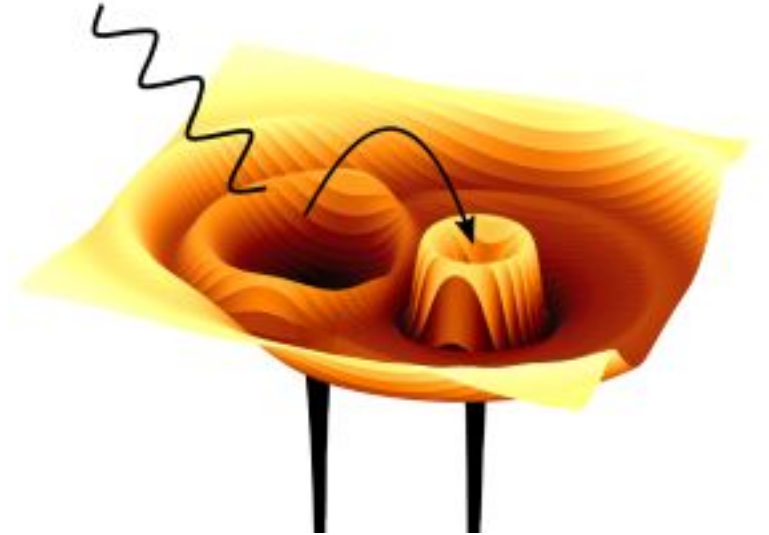
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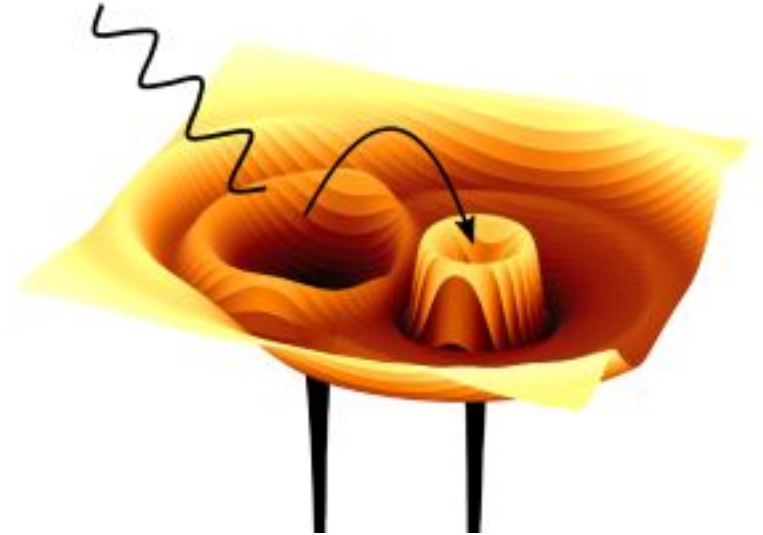


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This process can be modelled using:

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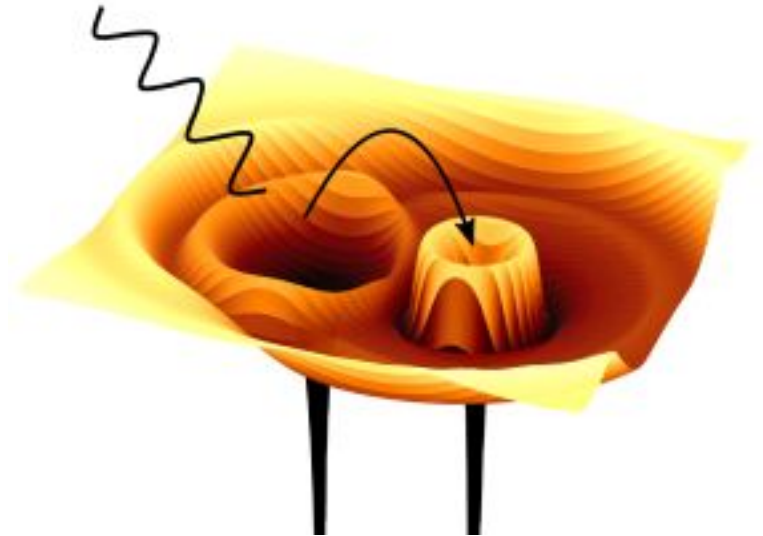


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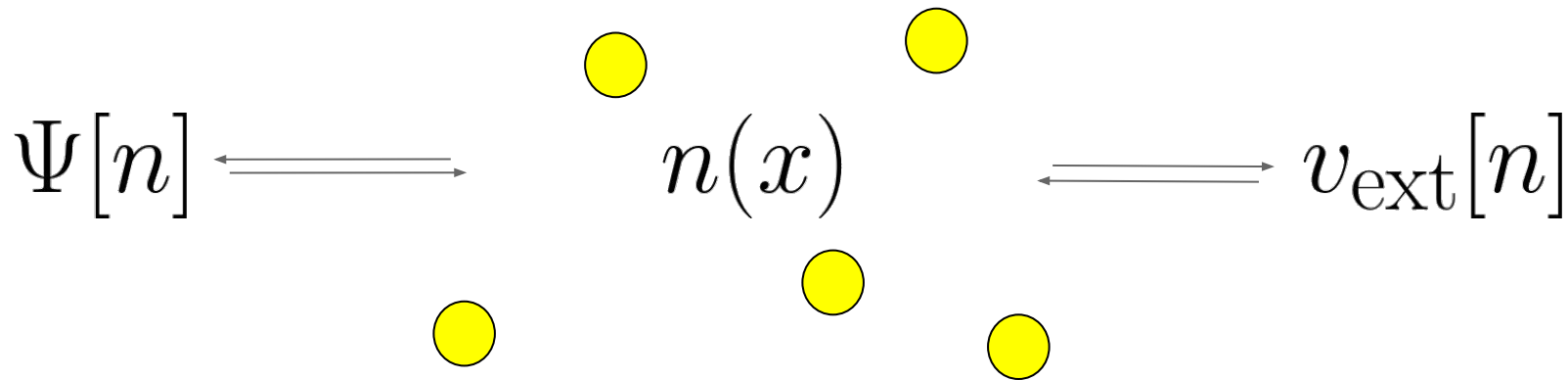
This process can be modelled using:

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- **Ensemble density functional theory**



# Ensemble density functional theory

Replace the **excited** wavefunction with the **excited** density –



E. K. Gross, L. N. Oliveira, and W. Kohn, Phys. Rev. A **37**, 2809 (1988)

E. K. Gross, L. N. Oliveira, and W. Kohn, Phys. Rev. A **37**, 2805 (1988)

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$$n(x) = \delta |\phi_{N+1}(x)|^2 + (1 - \delta) |\phi_N(x)|^2 + \sum_{i=1}^{N-1} |\phi_i(x)|^2$$

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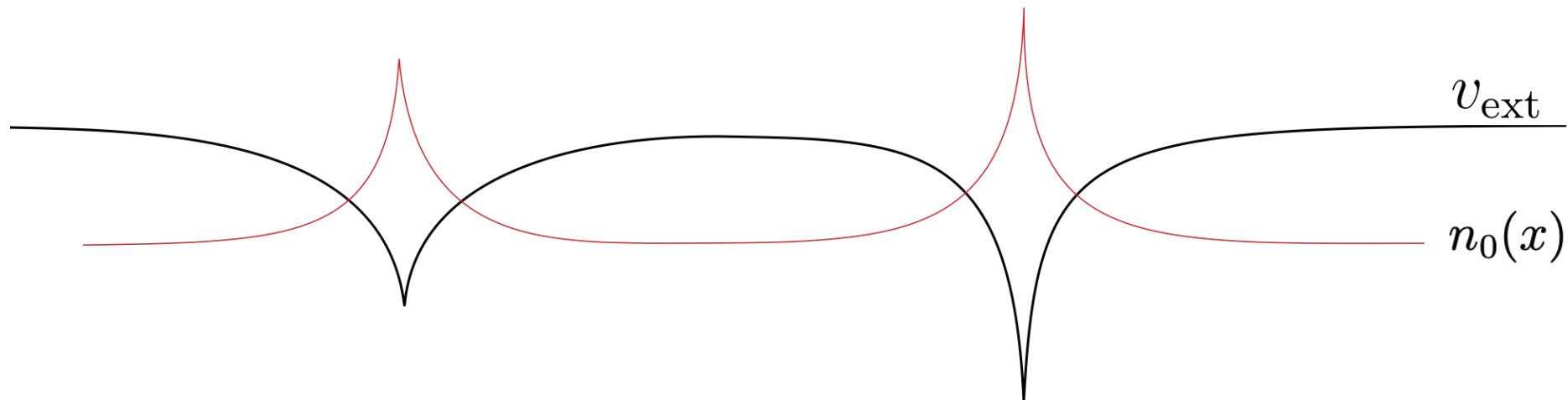
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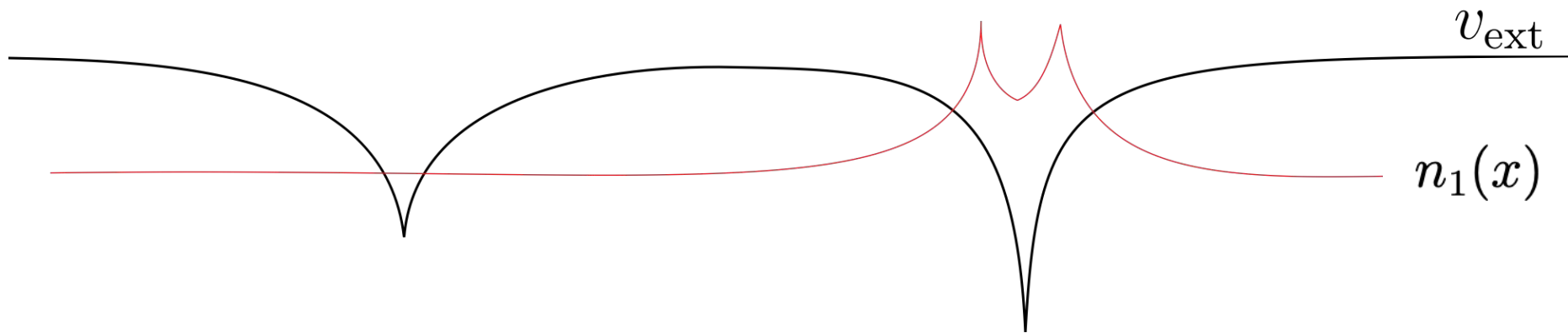
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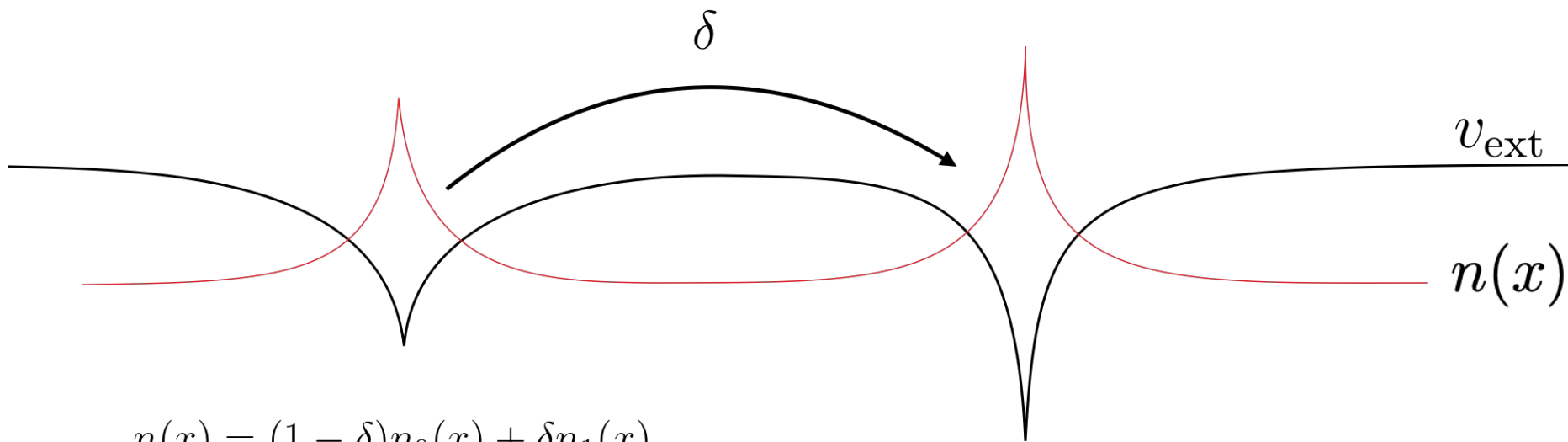
# Charge transfer excitation in a 1D molecule



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$$n(x) = (1 - \delta)n_0(x) + \delta n_1(x)$$

$$0 < \delta \ll 1$$

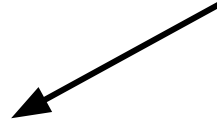
iDEA code

Define a 1/2/3 electron model  
system in 1D with a  
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Compute **exact** (time-dependent), fully-correlated  
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Reverse engineer **exact** density to give **exact**  
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Compute **approximate** (time-dependent) density  
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Reverse engineer **exact** density to give **exact**  
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Available approximations:



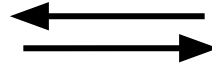
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- DFT and TDDFT – (generalised) Kohn-Sham (LDA, some novel methods)
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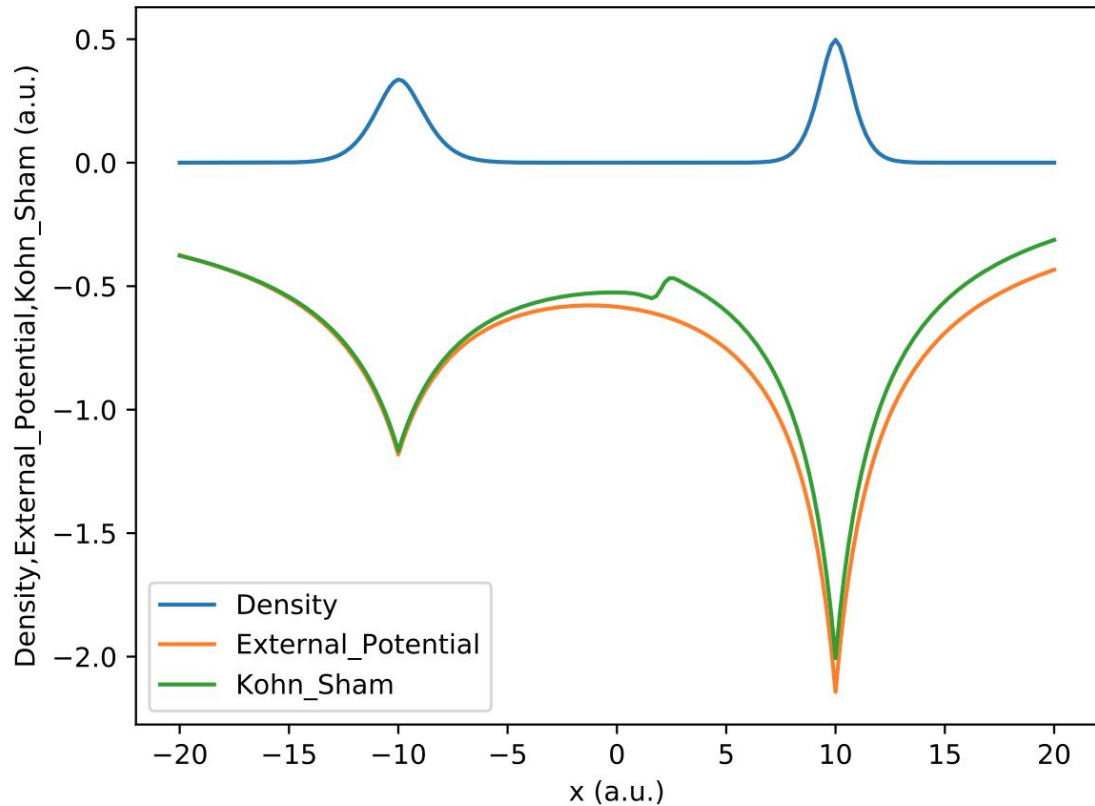
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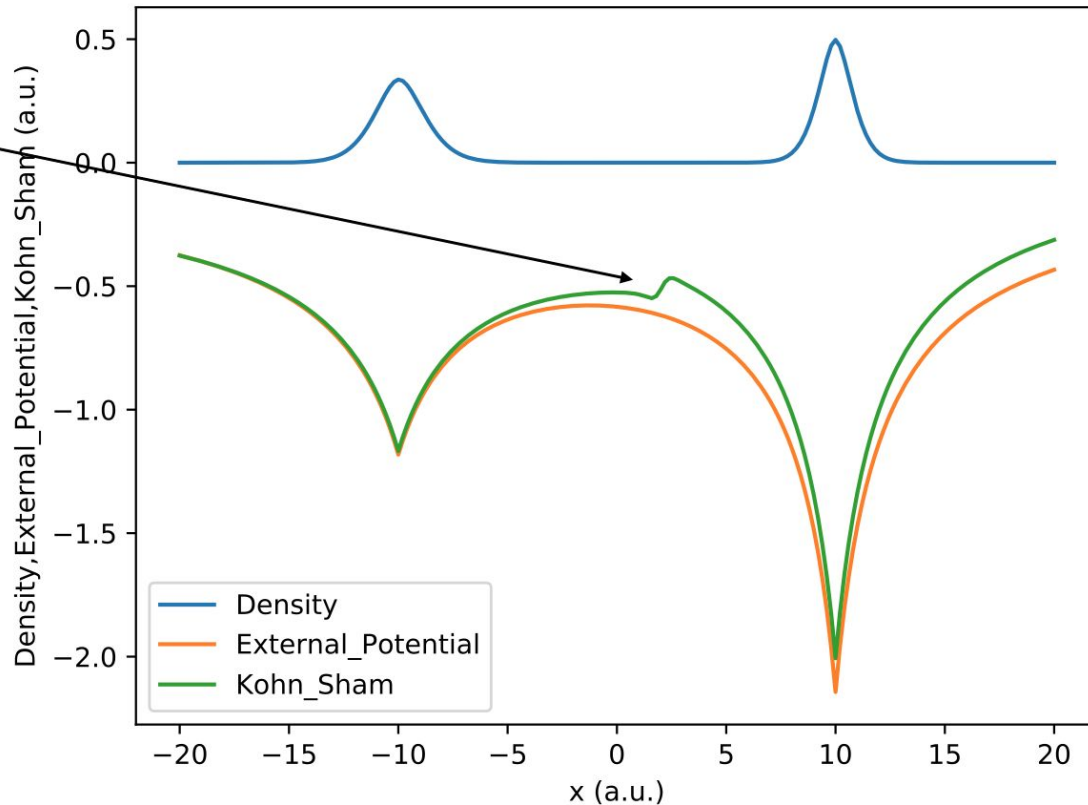
# Ground-state 1D molecule





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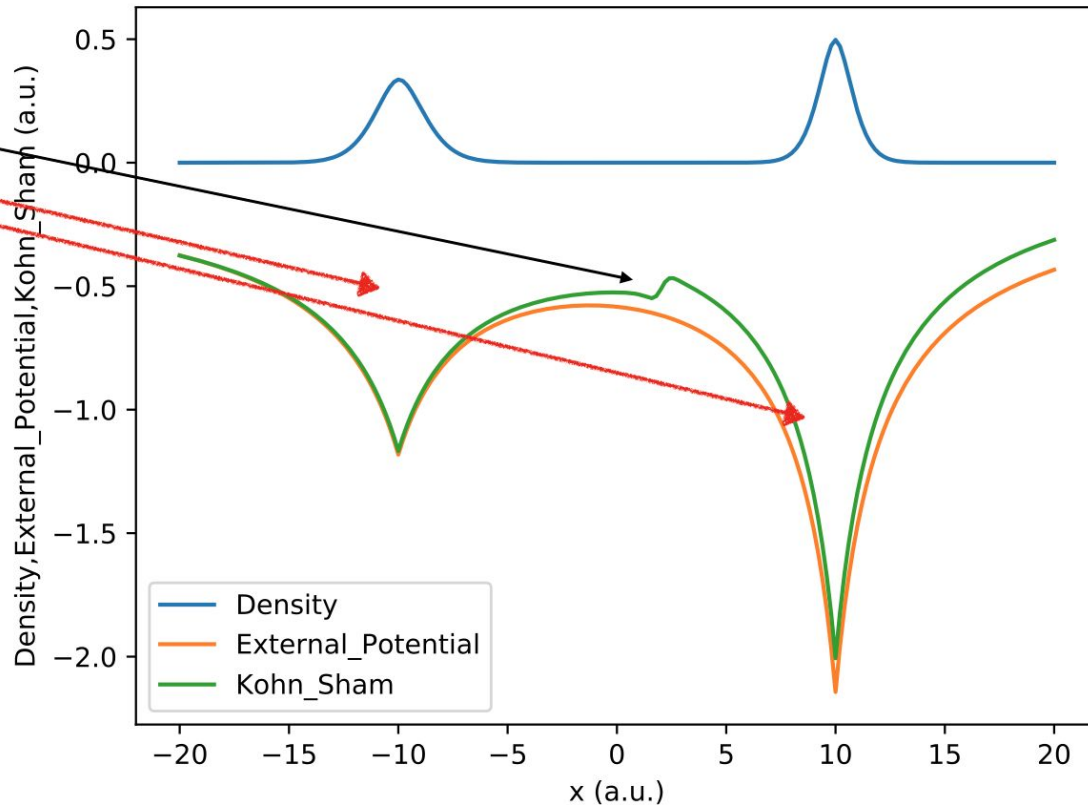
$$S = I_R - I_L + \eta_R^{\text{ho}} - \eta_L^{\text{ho}}$$



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Ionisation energy

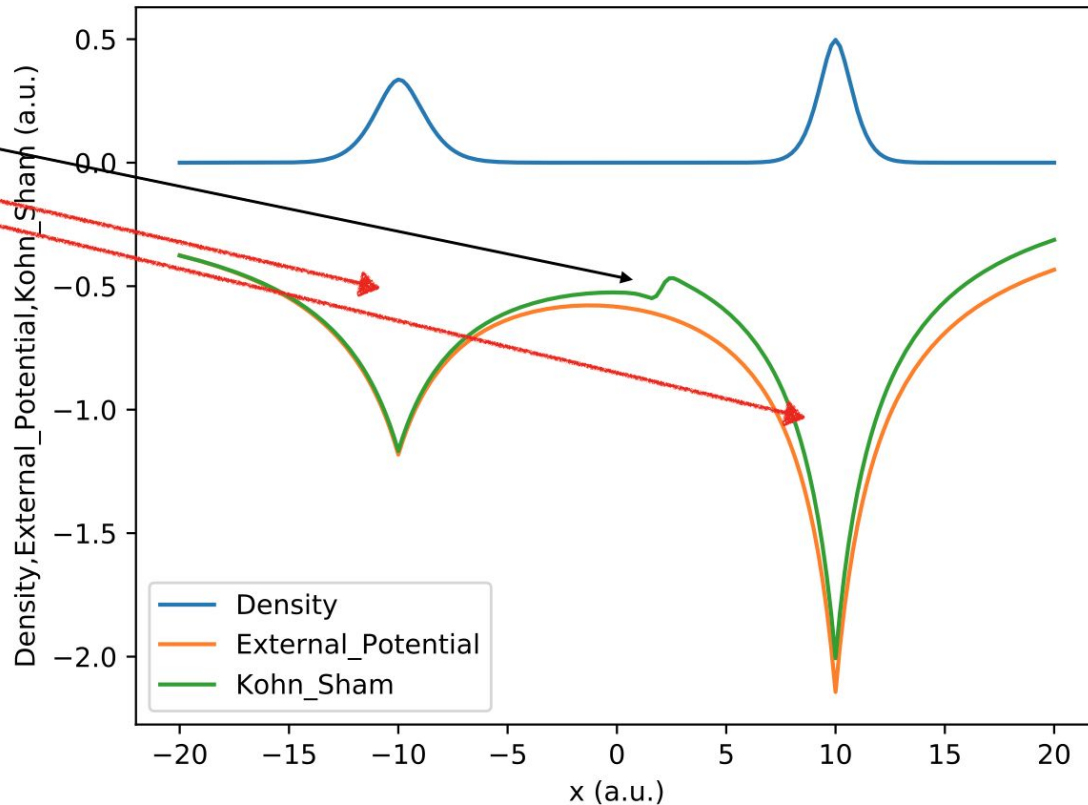


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Kohn-Sham molecular energies



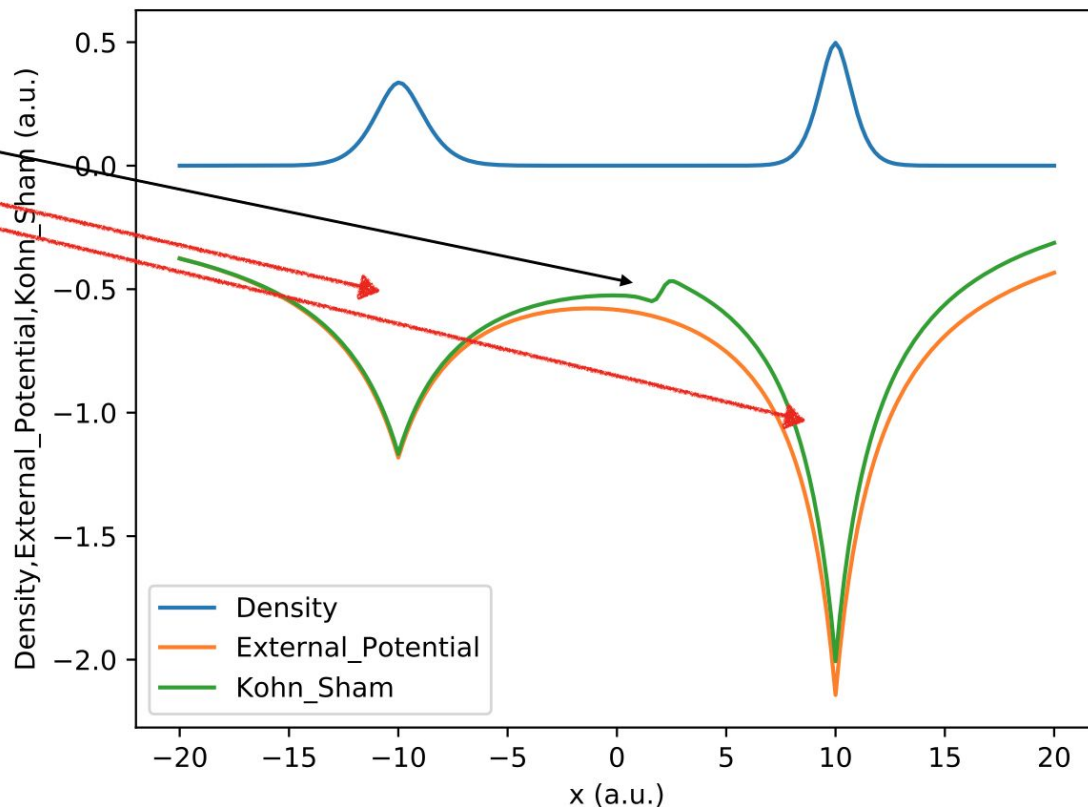
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Bonded case:  $\eta_R^{\text{ho}} = \eta_L^{\text{ho}}$



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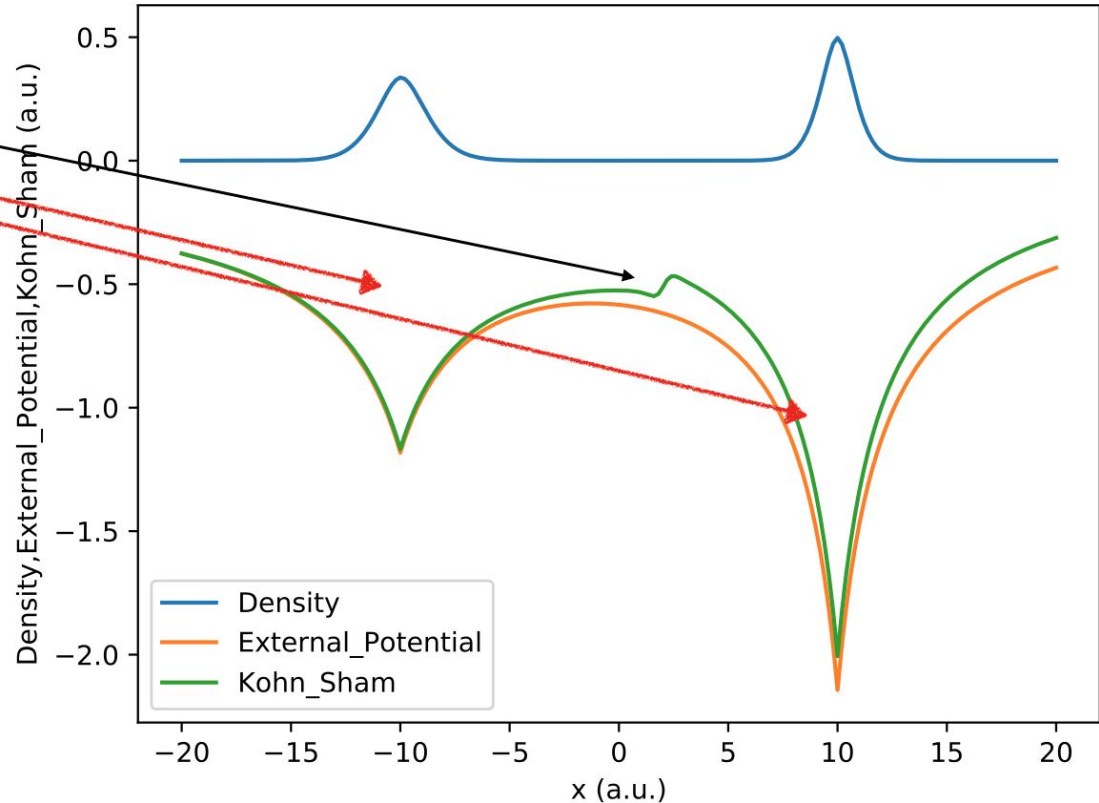
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$$\Rightarrow S = I_R - I_L$$



Almladh, C. O.; von Barth, U. In *Density Functional Methods in Physics*; Dreizler, R. M., da Providência, J., Eds.; NATO ASI Series; Plenum Press, 1985; Vol. 123, pp 209–231.

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The 'local number of electrons' of the right atom will increase by  $\delta$

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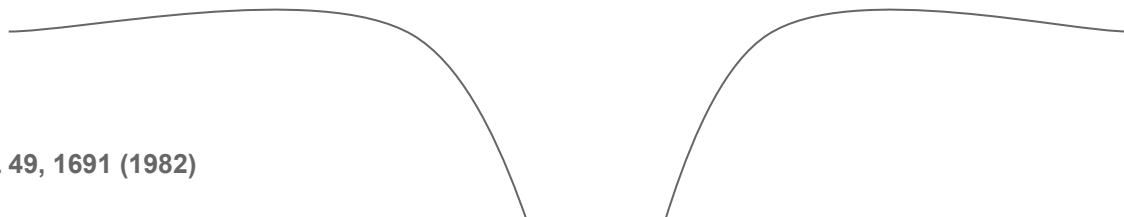
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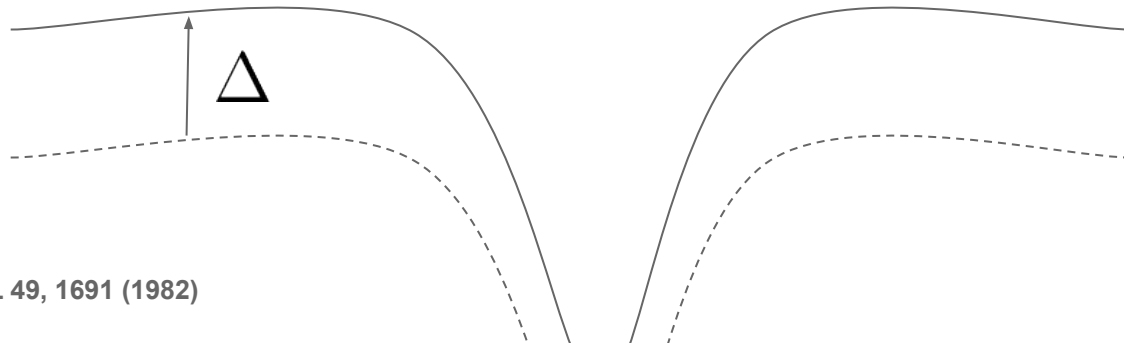


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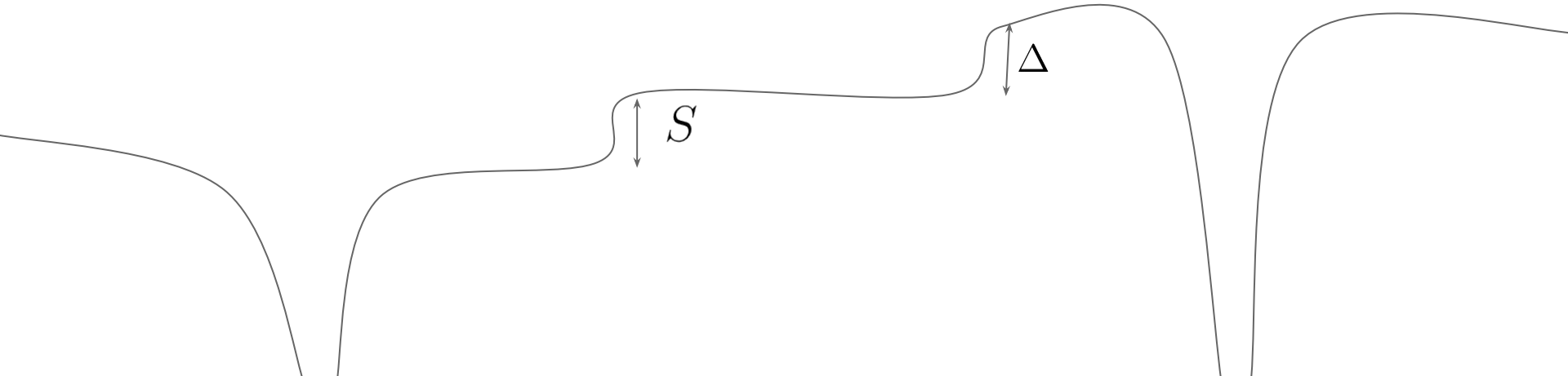
$$\Delta = I - A - (\varepsilon^{\text{lu}} - \varepsilon^{\text{ho}})$$



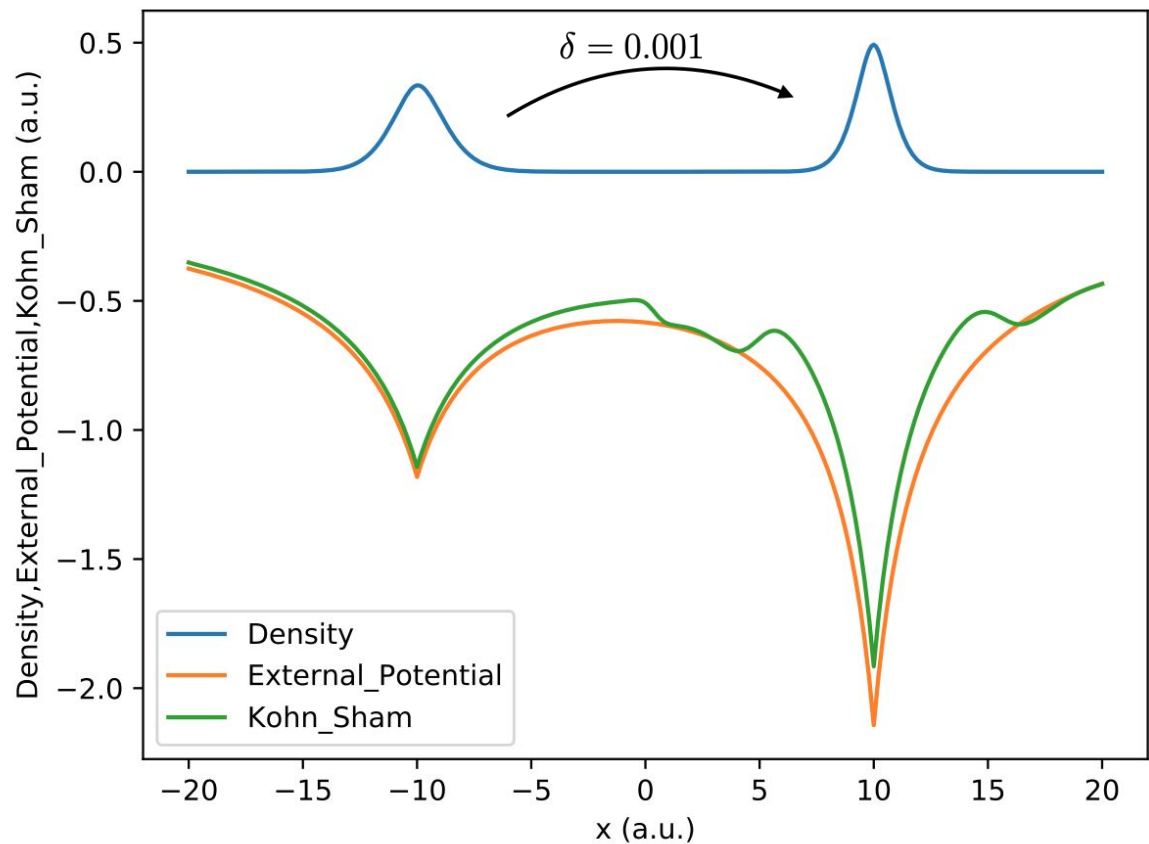
# When charge is transferred what do we expect?

What about the interatomic step,  $S$ ?

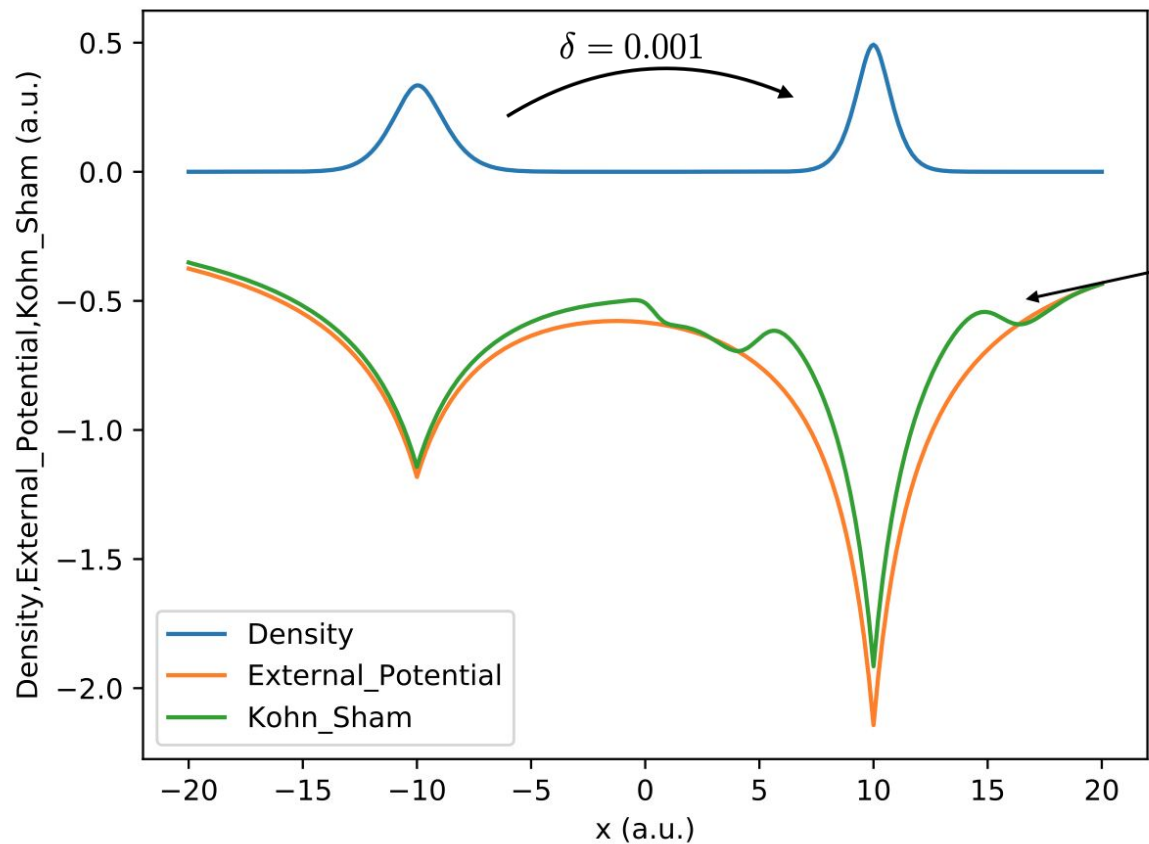
If  $S \rightarrow S + \Delta$  then the density will not be distributed correctly through the molecule!



# Charge transfer excitation in a 1D molecule



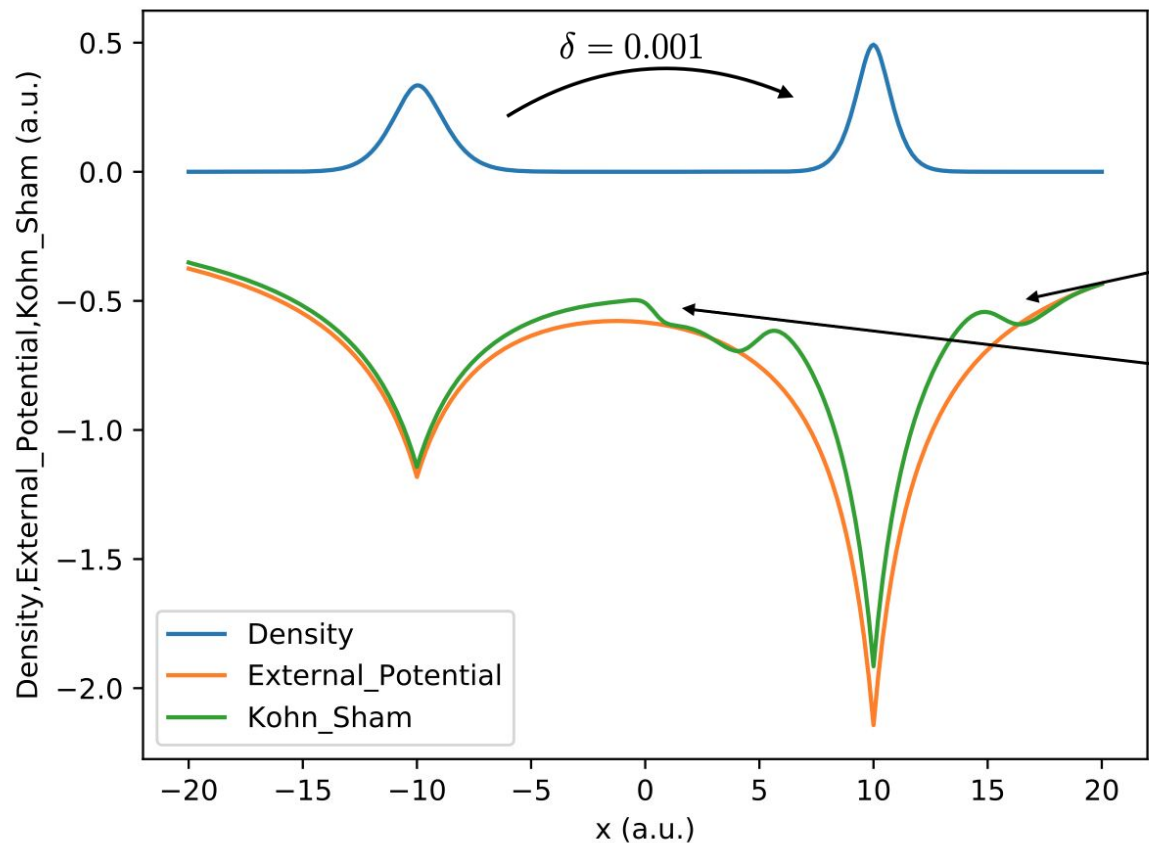
# Charge transfer excitation in a 1D molecule



Derivative discontinuity of Atom R:

$$\Delta_R = I_R - A_R + \varepsilon_R^{\text{ho}} - \varepsilon_R^{\text{lu}}$$

# Charge transfer excitation in a 1D molecule



Derivative discontinuity of Atom R:

$$\Delta_R = I_R - A_R + \varepsilon_R^{\text{ho}} - \varepsilon_R^{\text{lu}}$$

$$\Delta_{L \rightarrow R}^{\text{CT}} = I_L - A_R + \eta_L^{\text{ho}} - \eta_R^{\text{lu}}$$

What's this?

# Charge-transfer derivative discontinuity

$$\Delta_{L \rightarrow R}^{\text{CT}} = I_L - A_R + \eta_L^{\text{ho}} - \eta_R^{\text{lu}}$$

Derivative discontinuity of Atom R:

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# Charge-transfer derivative discontinuity

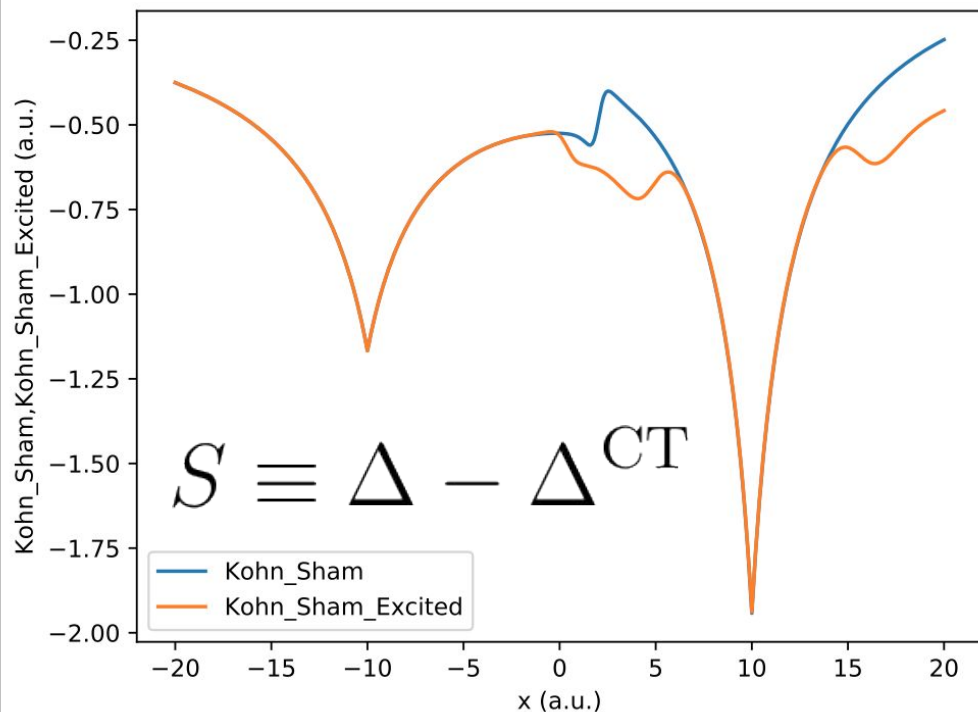
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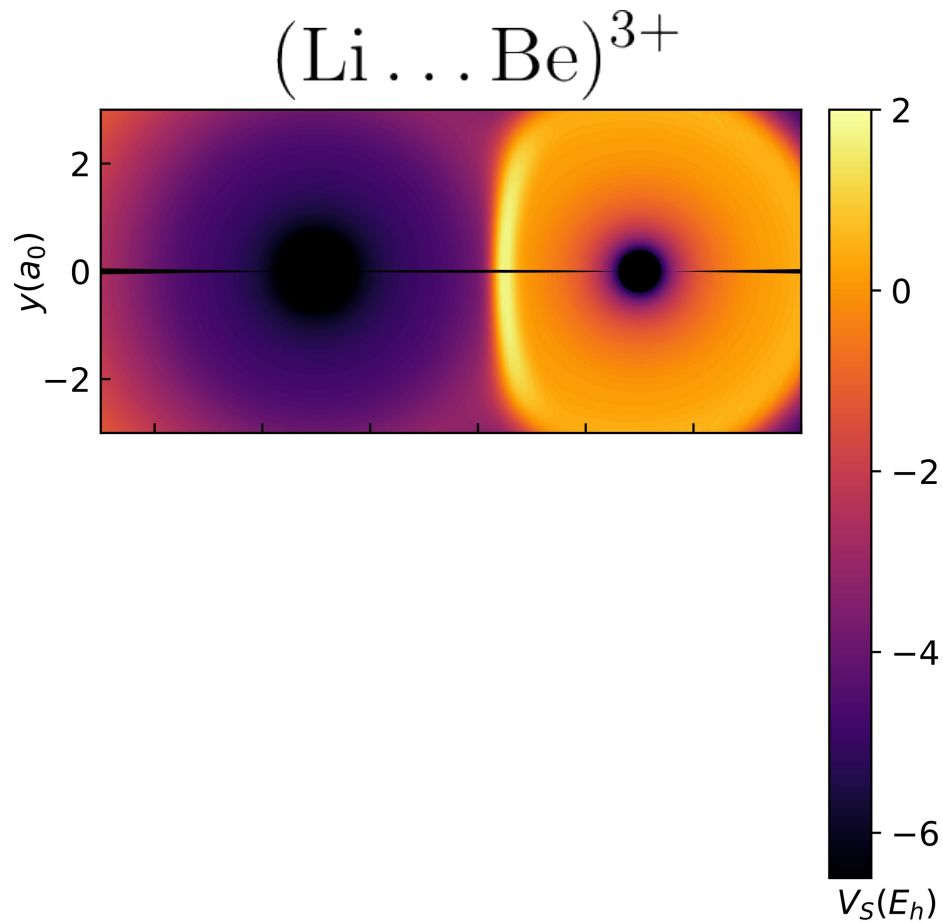
$$S = S_1 + S_2 = \Delta_R - \Delta_{L \rightarrow R}^{\text{CT}} = I_R - I_L + \eta_R^{\text{ho}} - \eta_L^{\text{ho}}$$

# Charge-transfer derivative discontinuity



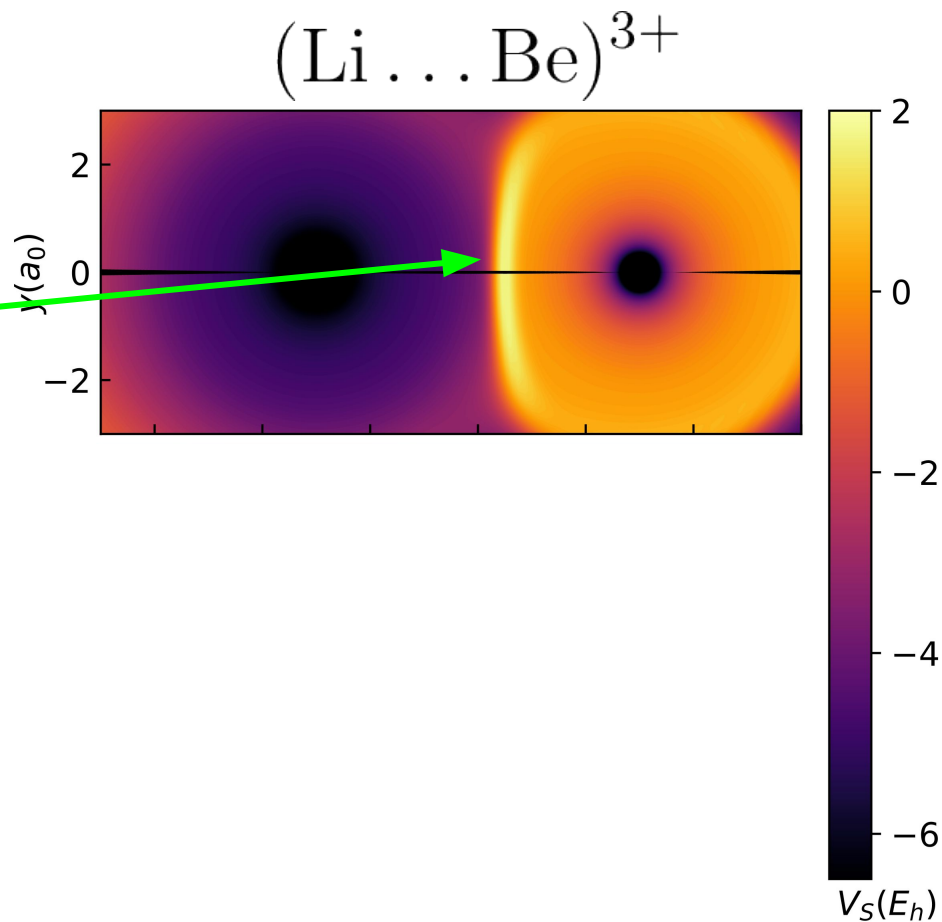


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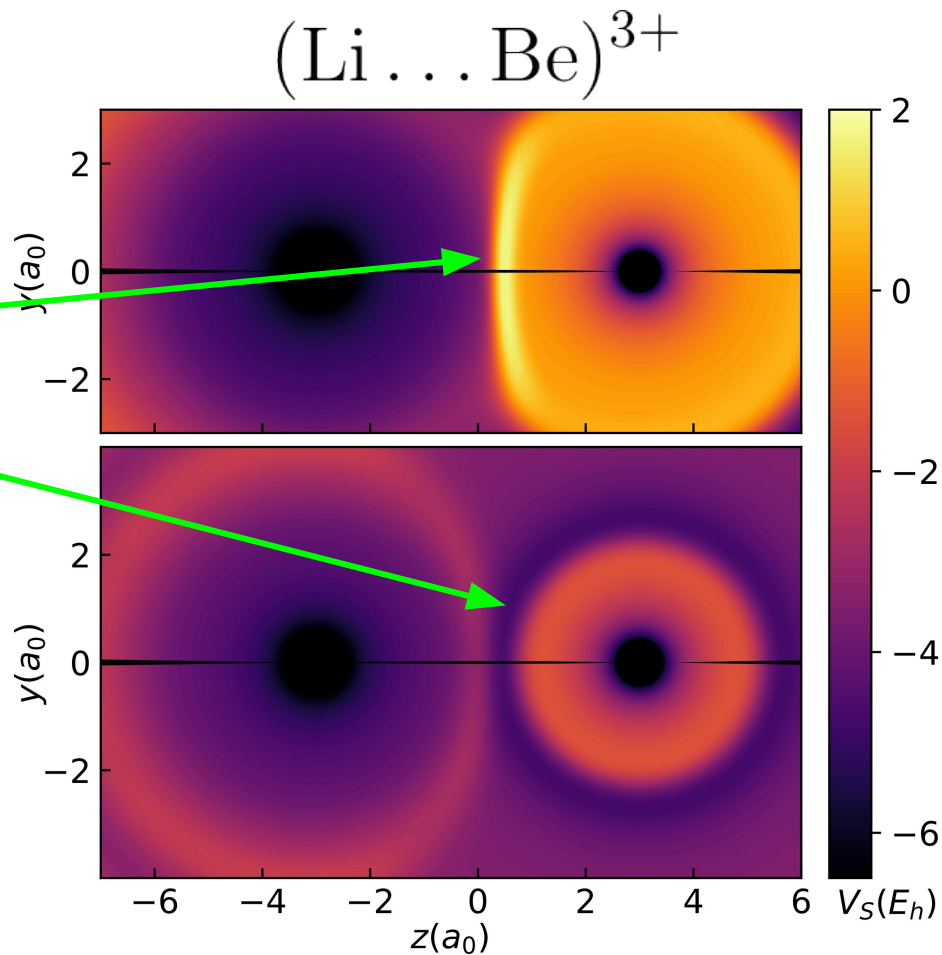
$$S = I_{\text{Be}} - I_{\text{Li}} + \eta_{\text{Be}}^{\text{ho}} - \eta_{\text{Li}}^{\text{ho}}$$



# What about in 3D?

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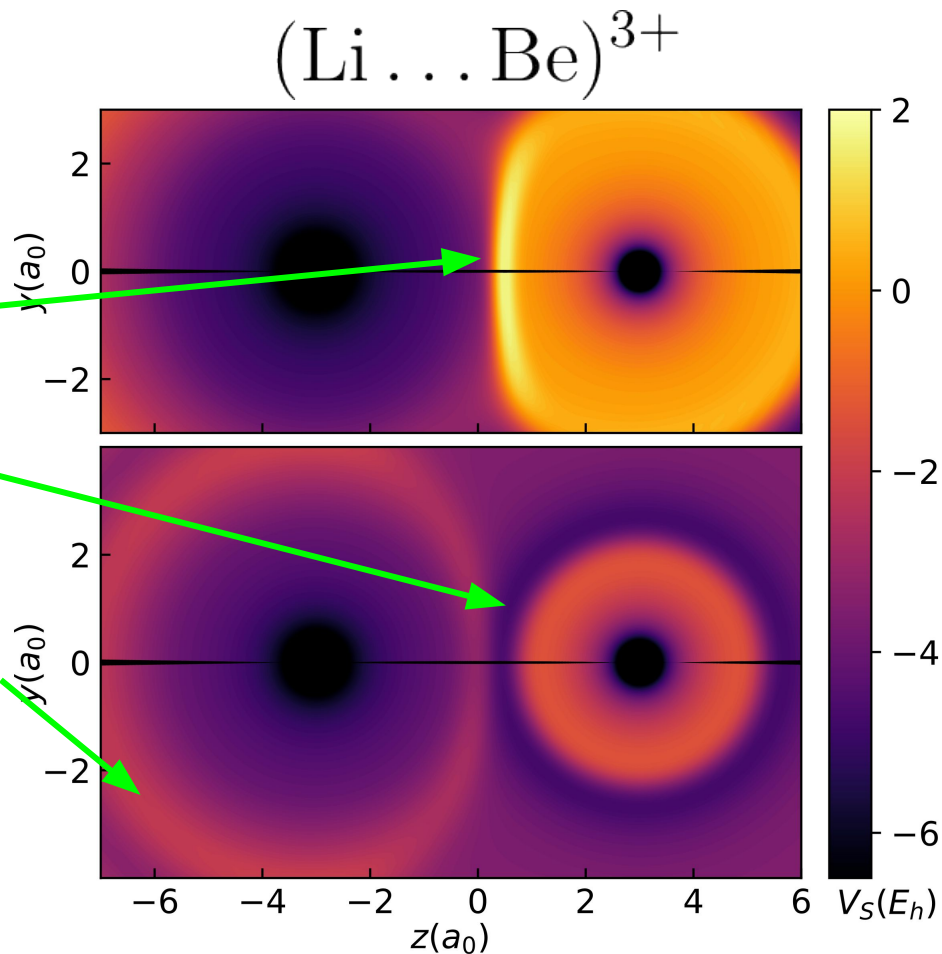


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$$\Delta_{\text{Li} \rightarrow \text{Be}}^{\text{CT}} = I_{\text{Li}} - A_{\text{Be}} + \eta_{\text{Li}}^{\text{ho}} - \eta_{\text{Be}}^{\text{lu}}$$



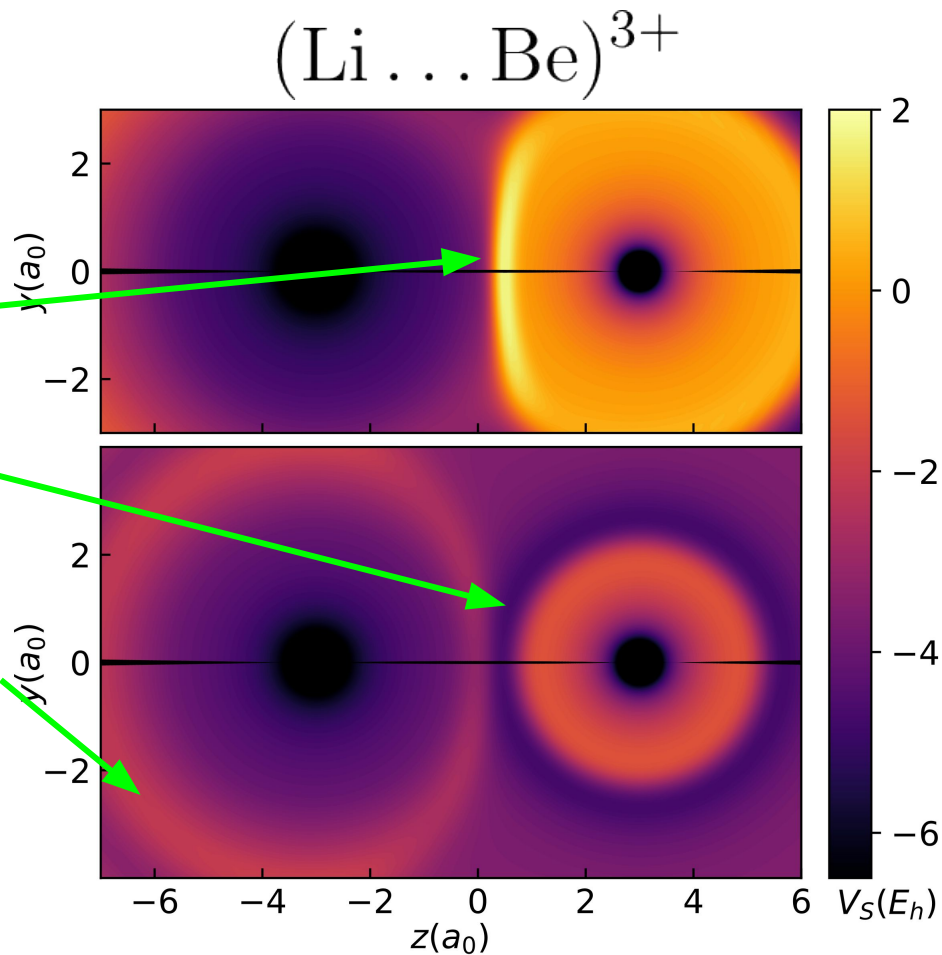
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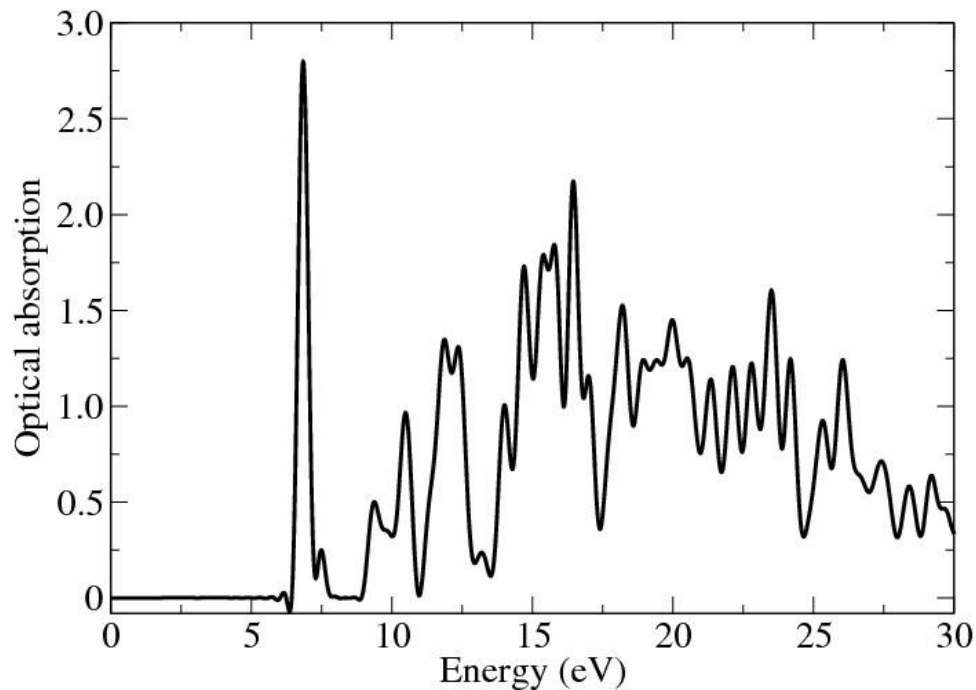
$$\Delta_{\text{Be}} = I_{\text{Be}} - A_{\text{Be}} + \eta_{\text{Be}}^{\text{ho}} - \eta_{\text{Be}}^{\text{lu}}$$

$$\Delta_{\text{Li} \rightarrow \text{Be}}^{\text{CT}} = I_{\text{Li}} - A_{\text{Be}} + \eta_{\text{Li}}^{\text{ho}} - \eta_{\text{Be}}^{\text{lu}}$$

$$S \equiv \Delta_{\text{Be}} - \Delta_{\text{Li} \rightarrow \text{Be}}^{\text{CT}}$$



# What about time-dependent density functional theory?



$$f_{xc}(x, x', t - t') = \frac{\delta v_{xc}(x, t)}{\delta n(x', t')}$$

# Conclusions

1. Charge transfer in molecules can be modelled via ensemble density functional theory
2. Discontinuities in the exchange-correlation energy lead to step features in the exact Kohn-Sham potential
3. Approximate density-functionals must capture these nonlocal features in order to yield accurate charge-transfer energies and densities (even in the ground state)

**M. J. P. Hodgson\***, **Eli Kraisler\***, **Axel Schild**, and **E. K. U. Gross**, *The Journal of Physical Chemistry Letters* 2017, 8 (24), 5974–5980.

My webpage: [www-users.york.ac.uk/~mjph501](http://www-users.york.ac.uk/~mjph501)

My email: [matthew.hodgson@mpi-halle.mpg.de](mailto:matthew.hodgson@mpi-halle.mpg.de)