

Dissipation and Ideal MHD

Contents

- Limits of ideal MHD
- Missing physics
- Layer models

Assuming collision dominated, so close to Maxwellian

- neglects the effect of that dissipation
- Fusion plasmas are nearly collisionless

Empirically Ideal MHD is found to work well

Because

- 1) The parallel dynamics which the model gets wrong are usually not important for equilibrium and large-scale instabilities
- 2) The perpendicular force balance is captured quite accurately by ideal MHD

Some missing physics:

- 1) Viscosity, due to collisions or Finite Larmor Radius
- 2) Resistivity
- 3) Wave-particle resonances
 - Landau closure (Landau damping)

Viscosity

Ion momentum

$$\text{Tr}[\underline{\underline{\Pi}}] = 0$$

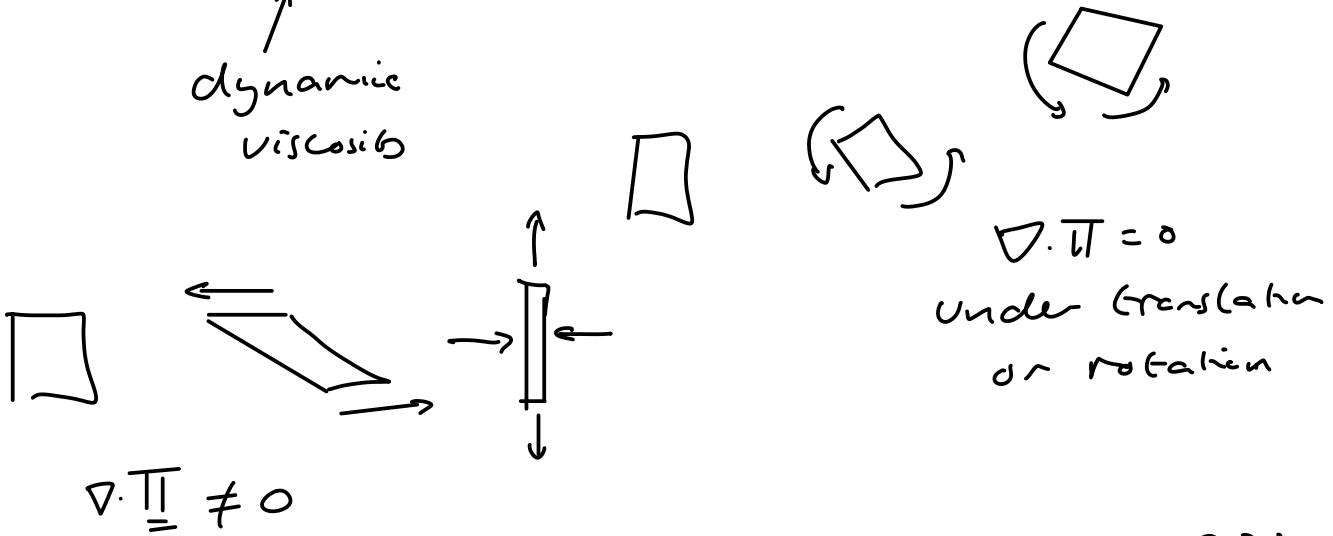


$$\rho \left[\frac{\partial \underline{u}}{\partial t} + \underline{u} \cdot \nabla \underline{u} \right] = \underline{j} \times \underline{B} - \nabla p - \underline{\underline{\nabla \cdot \underline{\Pi}}}$$

Stress tensor
(viscosity)

$$-\underline{\underline{\nabla \cdot \underline{\Pi}}} = \mu \nabla^2 \underline{u} + \frac{1}{3} \mu \nabla (\nabla \cdot \underline{u})$$

↑
dynamic
viscosity



$$\mu \text{ dynamic} = \rho \nu$$

$$\nu \text{ kinematic } [m^2/s]$$

Normalise u_0 speed
 L length
 L/u_0 time

$$\hat{\underline{u}} = \underline{u}/u_0 \quad \hat{\nabla} = \nabla L$$

$$\rho \frac{u_0^2}{L} \left(\frac{\partial \hat{u}}{\partial \hat{t}} + \hat{u} \cdot \hat{\nabla} \hat{u} \right) = \dots \underbrace{\mu \frac{u_0}{L^2}}_{\rho \nu} \left(\hat{\nabla}^2 \hat{u} + \frac{1}{3} \hat{\nabla} (\hat{\sigma} \cdot \hat{u}) \right)$$

$$\frac{\partial \hat{u}}{\partial \hat{t}} + \hat{u} \cdot \hat{\nabla} \hat{u} = \dots + \underbrace{\frac{\nu}{u_0 L}}_{\frac{1}{Re}} \left(\hat{\nabla}^2 \hat{u} + \frac{1}{3} \hat{\nabla} (\hat{\sigma} \cdot \hat{u}) \right)$$

$$Re = \frac{u_0 L}{\nu}$$

$$Re = \frac{\text{viscous diffusion time}}{\text{advection time}}$$

Resistivity

$$\underline{E} + \underline{u} \times \underline{B} = \eta \underline{j}$$

note: $\eta_{\perp} \approx 1.96 \eta_{\parallel}$

$$\frac{\partial \underline{B}}{\partial \hat{t}} = \nabla \times (\underline{u} \times \underline{B}) - \underbrace{\nabla \times (\eta \underline{j})}_{\approx -\frac{\eta}{\mu_0} \nabla^2 \underline{B}} \quad \underline{j} = \frac{1}{\mu_0} \nabla \times \underline{B}$$

$$\approx -\frac{\eta}{\mu_0} \nabla^2 \underline{B}$$

Diffusion coefficient

Typical length L
 Speed u_0
 Magnetic field B_0

$$\underbrace{\frac{B_0 u_0}{L}} \frac{\partial \hat{B}}{\partial \hat{t}} = \underbrace{\frac{B_0 u_0}{L}} \hat{\nabla} \times (\hat{u} \times \hat{B}) + \frac{\eta}{\mu_0} \frac{1}{L^2} B_0 \hat{\nabla}^2 \hat{B}$$

$$\frac{\partial \underline{\hat{B}}}{\partial \hat{t}} = \underline{\hat{v}} \times (\underline{\hat{u}} \times \underline{\hat{B}}) + \frac{\eta}{\mu_0 u_0 L} \nabla^2 \underline{\hat{B}}$$

$$\frac{1}{R_m}$$

Magnetic Reynolds number

$$R_m = \frac{\text{Resistive diffusion timescale}}{\text{advection timescale}}$$

Set $u_0 = v_A$ Alfvén speed

$$S = \frac{\mu_0 v_A L}{\eta}$$

Lundquist number

Small scales and Layers

- if $Re \sim 1$ viscosity \sim advection

$$\frac{u_0 L}{\nu} \sim 1$$

$$\underline{L_\nu} \sim \frac{\nu}{u_0}$$

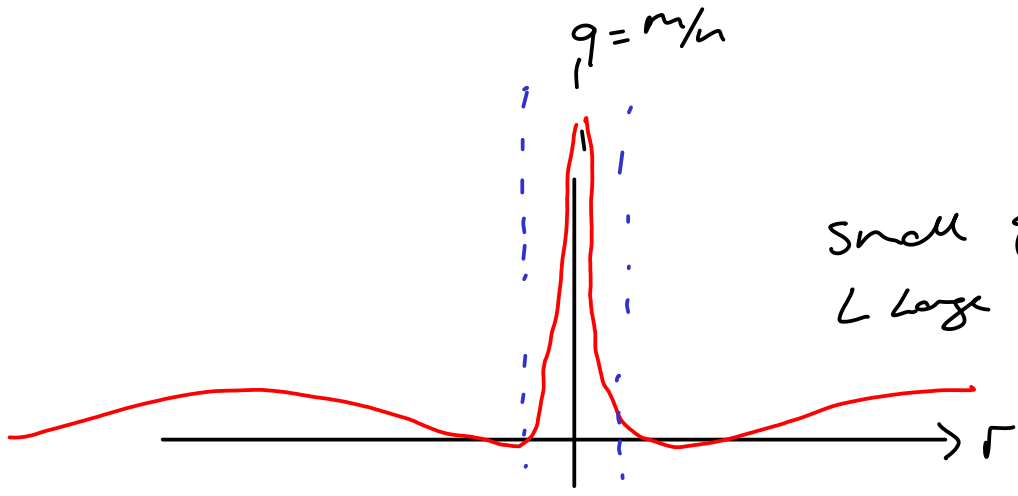
- if $R_m \sim 1$ resistive \sim advection

$$\frac{\mu_0 L u_0}{\eta} \sim 1$$

$$\underline{L_\eta} \sim \frac{\eta}{\mu_0 u_0}$$

if $L \gg L_\nu, L_\eta \Rightarrow$ dissipation is small
 \sim ideal MHD

if $L \sim L_\nu, L_\eta \Rightarrow$ non-ideal effects important



Small gradients
 L large \rightarrow ideal

small L
Non-ideal
Case