

# Poisson brackets in Reduced MHD

## Contents

- In field-aligned coordinates
- Conservation properties
- Numerical conservation laws
- Arakawa brackets

$$\underline{V}_E \cdot \nabla f = \frac{1}{B} \underline{b} \times \nabla \phi \cdot \nabla f$$

in field-aligned coordinates  $\underline{B} = \nabla x \times \nabla y$

$$(\underline{b} \times \nabla \phi) \cdot \nabla f = (\nabla \phi \times \nabla f) \cdot \underline{b}$$

$$\nabla \phi = \nabla x \frac{\partial \phi}{\partial x} + \nabla y \frac{\partial \phi}{\partial y} + \nabla z \frac{\partial \phi}{\partial z}$$

$$\begin{aligned} B \underline{V}_E \cdot \nabla f &= \left( \frac{\partial \phi}{\partial x} \frac{\partial f}{\partial y} - \frac{\partial \phi}{\partial y} \frac{\partial f}{\partial x} \right) \underbrace{(\nabla x \times \nabla y)}_B \cdot \underline{b} \\ &+ \left( \frac{\partial \phi}{\partial x} \frac{\partial f}{\partial z} - \frac{\partial \phi}{\partial z} \frac{\partial f}{\partial x} \right) \underbrace{(\nabla x \times \nabla z)}_{\frac{1}{J} \underline{e}_y} \cdot \underline{b} \\ &+ \left( \frac{\partial \phi}{\partial y} \frac{\partial f}{\partial z} - \frac{\partial \phi}{\partial z} \frac{\partial f}{\partial y} \right) \underbrace{(\nabla y \times \nabla z)}_{\frac{1}{J} \underline{e}_x} \cdot \underline{b} \end{aligned}$$

$$\underline{V}_E \cdot \nabla f \approx \frac{\partial \phi}{\partial x} \frac{\partial f}{\partial y} - \frac{\partial \phi}{\partial y} \frac{\partial f}{\partial x} \quad \begin{array}{l} \text{Coordinates} \\ \sim \text{orthogonal} \end{array}$$

$$= [\phi, f] \quad \text{Poisson bracket}$$

antisymmetric operator perpendicular to  $\underline{\beta}$

$$\frac{\partial n}{\partial t} + [\phi, n] = 0 \quad \text{incompressible flow}$$

$$\frac{1}{\beta} \underline{k} \times \nabla \phi \cdot \nabla f = \nabla \cdot \left( \frac{\underline{k} \times \nabla \phi}{\beta} f \right) - f \nabla \cdot \left( \frac{\underline{k} \times \nabla \phi}{\beta} \right)$$

$$\frac{\partial}{\partial t} \int n \, dV + \int \nabla \cdot \left( \frac{\underline{k} \times \nabla \phi}{\beta} f \right) \, dV = 0$$

conservation

$$\int [\phi, n] \, dV = 0 \quad \text{if no boundary flux}$$

numerically  $\sum_i [\phi, n]_i \Delta V_i = 0$

discrete form of conservation law

$$\frac{\partial \omega}{\partial t} + [\omega, \phi] = 0 \quad \omega = \nabla^2 \phi$$

$$\omega \frac{\partial \omega}{\partial t} + \omega [\omega, \phi] = 0 \quad (\text{incompressible})$$

$$\frac{1}{2} \frac{\partial}{\partial t} (\omega^2) \quad \downarrow \quad \frac{1}{2} \nabla \cdot \left( \omega^2 \frac{\underline{k} \times \nabla \phi}{\beta} \right)$$

$\Rightarrow \omega^2$  conserved (enstrophy)

2D grid  $i, j$

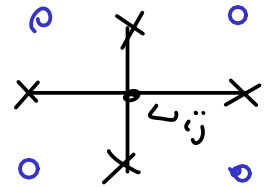
$$\frac{\partial}{\partial t}(\omega_{ij}^2) + \sum_{i,j} \omega_{ij} [\phi, \omega]_{ij} = 0 \quad \text{discrete form}$$

should = 0 if  $\frac{\partial}{\partial t} = 0$

2<sup>nd</sup> order central difference

$$\frac{\partial \phi}{\partial x} \Big|_{ij} \approx \frac{\phi_{i+1,j} - \phi_{i-1,j}}{2\Delta x}$$

$$[\phi, \omega] \approx \frac{(\phi_{i+1,j} - \phi_{i-1,j})(\omega_{i,j+1} - \omega_{i,j-1})}{4\Delta x \Delta y} - \frac{(\phi_{i,j+1} - \phi_{i,j-1})(\omega_{i+1,j} - \omega_{i-1,j})}{4\Delta x \Delta y}$$



$$\omega_{ij} [\phi, \omega]_{ij} = \omega_{ij} (\phi_{i+1,j} - \phi_{i-1,j}) \omega_{i,j+1} + \dots$$

$$\omega_{i,j+1} [\phi, \omega]_{i,j+1} = \omega_{i,j+1} (\phi_{i+1,j+1} - \phi_{i-1,j+1}) (-\omega_{ij}) + \dots$$

⇒ Terms do not balance

⇒ breaks conservation of enstrophy

Arakawa (1966)

→ conserves total  $n$ ,  $\omega$  and  $\omega^2$