

7.1

Lecture 7

- Review types of flow

Incompressible $\nabla \cdot \underline{u} = 0$

Irrotational $\nabla \times \underline{u} = 0$

- Derivative along the flow
Momentum equation $\frac{D}{Dt} = \frac{\partial}{\partial t} + \underline{u} \cdot \nabla$

$$\frac{\partial \underline{u}}{\partial t} + (\underline{u} \cdot \nabla) \underline{u} = -\nabla \left(\frac{P}{\rho} + \chi \right)$$

$$\chi = gz \quad \text{gravitational potential}$$

$$\text{so } \underline{g} = -\nabla \chi$$

Identities: $(\underline{u} \cdot \nabla) \underline{u} = (\nabla \times \underline{u}) \times \underline{u} + \nabla \left(\frac{1}{2} u^2 \right)$

$$\Rightarrow \frac{\partial \underline{u}}{\partial t} + (\nabla \times \underline{u}) \times \underline{u} = -\nabla \left(\frac{P}{\rho} + \frac{1}{2} u^2 + \chi \right)$$

$\underbrace{\hspace{10em}}_H$

Steady flow

$$(\nabla \times \underline{u}) \times \underline{u} = -\nabla H$$

Dot with \underline{u} :

$$(\underline{u} \cdot \nabla) H = 0 \quad \text{i.e. } H \text{ constant along flow}$$

Bernoulli streamline theorem

$$\text{If } \nabla \times \underline{u} = 0 \text{ then } \nabla H = 0 \quad \text{i.e. } H \text{ constant}$$

$$\ast \text{ idea of vorticity } \nabla \times \underline{u} = \underline{\omega}$$

$\rho v^\alpha = \text{const}$

polytropic $\gamma = 3$

Adiabatic

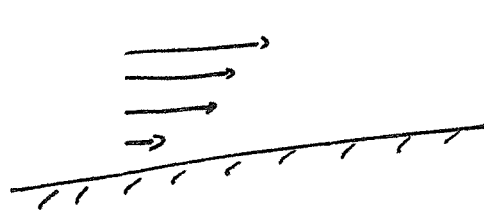
isobaric $\alpha = 0$

isothermal $\alpha = 1$

adiabatic $\alpha = \gamma$

Vorticity: local rotation of fluid elements

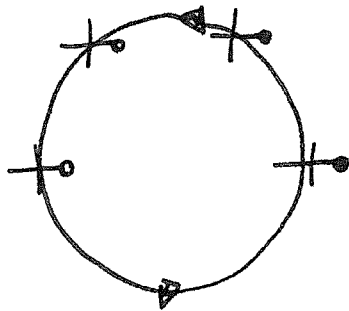
e.g.



$$\underline{u} = (\beta y, 0, 0)$$

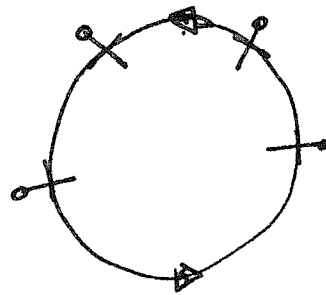
$$\Rightarrow \underline{\omega} = -\beta \hat{z}$$

no global rotation, but has vorticity



$$u_0 \propto \frac{1}{r}$$

no vorticity

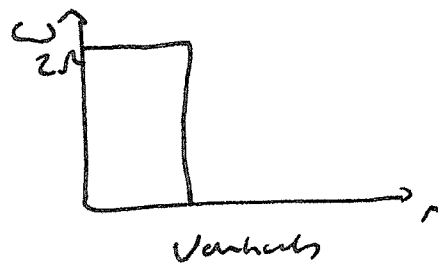
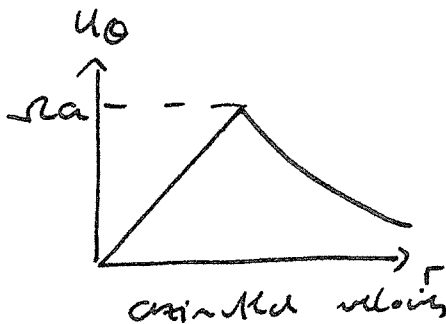


$$u_0 \propto r$$

has vorticity

Rankine Vortices

$$u_\theta = \begin{cases} \Omega r & r < a \\ \frac{\Omega a^2}{r} & r > a \end{cases}$$



Steady flow past a wing at small angles of incidence is typically irrotational

7.2

Proof: The vorticity equation

$$\frac{\partial \underline{u}}{\partial t} + \underline{\omega} \times \underline{u} = -\nabla H$$

Take curl

$$\frac{\partial \underline{\omega}}{\partial t} + \nabla \times (\underline{\omega} \times \underline{u}) = 0$$

NB: Pressure removed

Vector identities

$$\frac{\partial \underline{\omega}}{\partial t} + (\underline{u} \cdot \nabla) \underline{\omega} - (\underline{\omega} \cdot \nabla) \underline{u} + \underline{\omega} \nabla \cdot \underline{u} - \underline{u} \nabla \cdot \underline{\omega} = 0$$

Incompress. $\nabla \cdot \underline{u} = 0$

$$\Rightarrow \frac{\partial \underline{\omega}}{\partial t} + (\underline{u} \cdot \nabla) \underline{\omega} = (\underline{\omega} \cdot \nabla) \underline{u}$$

$$\boxed{\frac{D \underline{\omega}}{Dt} = (\underline{\omega} \cdot \nabla) \underline{u}}$$

In 2D flow

$$\underline{u} = [u_x(x, y, z, t), u_y(x, y, z, t), 0]$$

$$\underline{\omega} = [0, 0, \omega_z]$$

$$\text{So } (\underline{\omega} \cdot \nabla) \underline{u} = \omega_z \frac{\partial \underline{u}}{\partial z} = 0$$

and so $\boxed{\frac{D \underline{\omega}}{Dt} = 0}$

In 2D flow of an ideal incompressible fluid the vorticity of each fluid element is conserved

$$\text{steady flow } (\underline{u} \cdot \nabla) \omega_z = 0$$

If $\underline{\omega} = \nabla \times \underline{v} = 0$ everywhere then the flow is irrotational. In this case $\nabla H = 0$

and so

$$H = \frac{P}{\rho} + \frac{1}{2}u^2 + \chi$$

is a constant everywhere.

Irrotational flow allows the velocity to be written as a potential ϕ :

$$\underline{v} = \nabla \phi \quad (\text{so } \nabla \times \underline{v} = \nabla \times \nabla \phi = 0)$$

Greatly simplifies analysis, played significant role in development of aerodynamics.

However:

1. This flow is dissipationless (ideal)
So doesn't cover friction, drag forces
2. As we shall see, irrotational flow is not unique. It allows arbitrary discontinuities at the surface of a body
3. In 2D the flow is only unique if a condition is also provided by col-loc assumptions (e.g. Kutta-Joukowski condition)