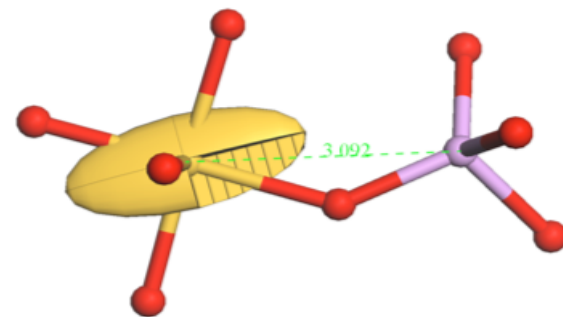
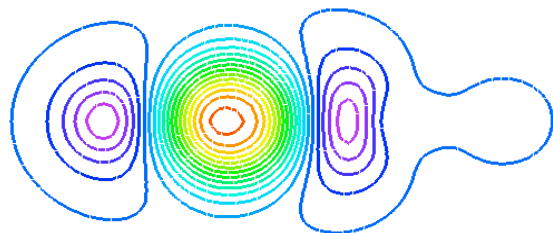


THE UNIVERSITY *of York*

Solid-state Nuclear Magnetic Resonance (NMR)



Phil Hasnip

Condensed Matter Dynamics Group

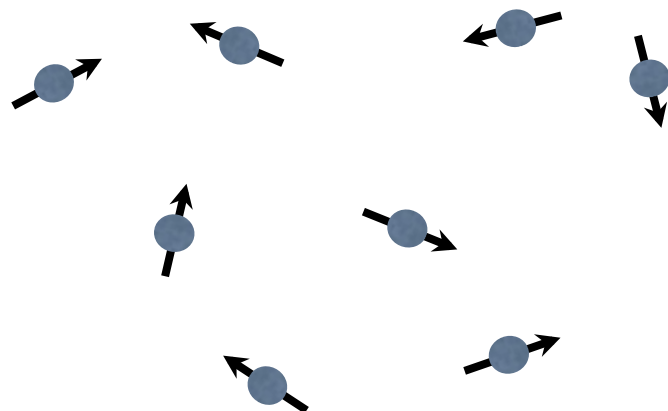
Department of Physics,

University of York, U.K.

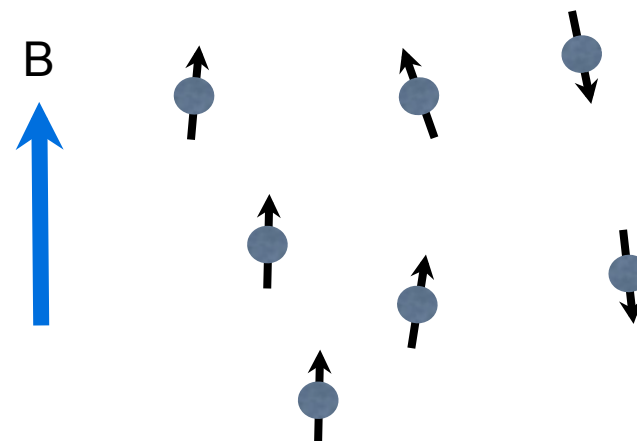
<http://www-users.york.ac.uk/~pjh503>

Many slides courtesy of Jonathan Yates (University of Oxford)

Zero Field



Applied Field

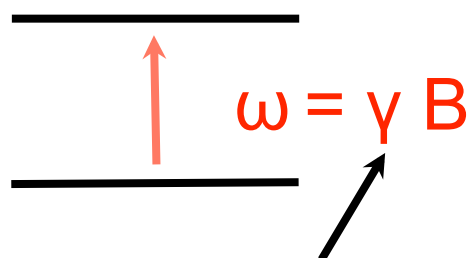


NMR signal is proportional to population difference
For ^1H in 9.4T at 298K

$$\Delta E = 2.65 \times 10^{-25} \text{ J}$$

$$n_{\text{upper}}/n_{\text{lower}} = \exp(-\Delta E / kT) = 0.999935$$

Sensitivity is an issue! Need Large γ and/or B



magnetogyric ratio
fundamental nuclear constant

Isotope	Spin	γ $10^7 \text{ T}^{-1} \text{ rad s}^{-1}$	Freq MHz @ 9.4T	Abundance %
^1H	1/2	26.75	400	100
^{13}C	1/2	6.73	100.6	1.1
^{29}Si	1/2	-5.32	79.6	4.7
^{31}P	1/2	-10.84	162.1	100

Isotope	Spin	γ $10^7 \text{ T}^{-1} \text{ rad s}^{-1}$	Freq MHz @ 9.4T	Abundance %
^1H	1/2	26.75	400	100
^{13}C	1/2	6.73	100.6	1.1
^{29}Si	1/2	-5.32	79.6	4.7
^{31}P	1/2	-10.84	162.1	100
^{17}O	5/2	-3.63	54.1	0.04
^{49}Ti	7/2	-1.51	22.6	5.41



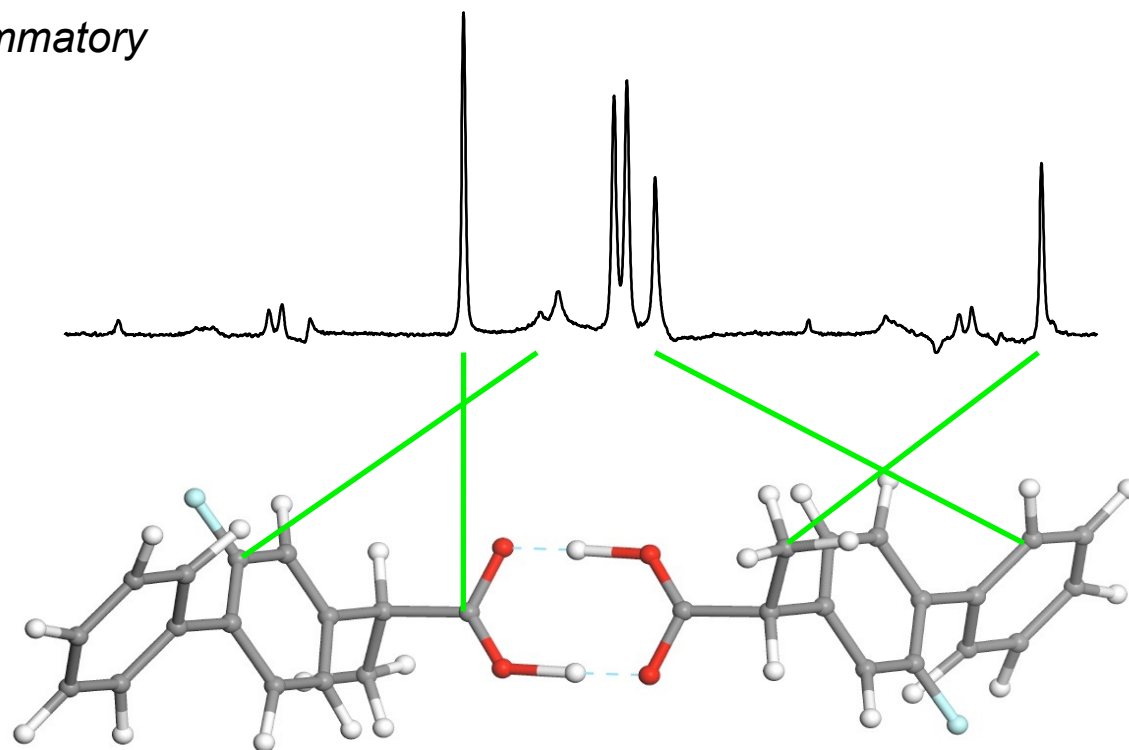
600 MHz 14.1 T
~£800,000

400 MHz 9.4 T
~£300,000

Image courtesy Sharon Ashbrook (St Andrews)

Flurbiprofen

non-steroidal anti-inflammatory



Each distinct C atom experiences a different magnetic field and resonates at a unique frequency.

Measure the change wrt a standard (for ^{13}C this is liquid tetramethylsilane)

*The local field a nucleus feels is not quite the same as the applied field (B_0). The local field is influenced by the positions of neighbouring atoms and the electronic structure of the material. Hence **NMR is a sensitive probe of atomic-level structure and dynamics.***

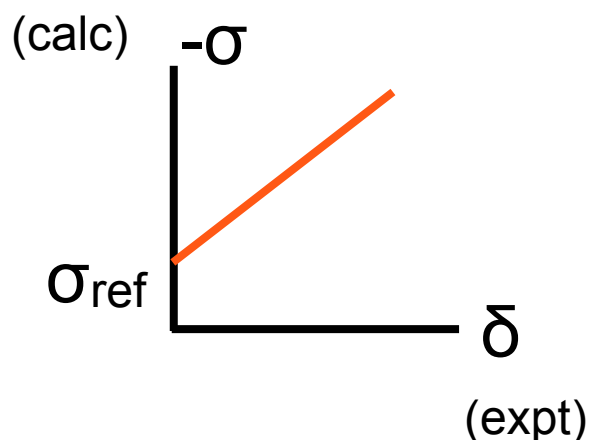
Larmor frequency $\omega = \gamma B$

σ magnetic shielding

$$B = (1 - \sigma) B_0$$

δ chemical shift

$$\delta = \sigma_{\text{ref}} - \sigma$$



σ_{ref} is the shielding of a nucleus in a “standard” material. For ^{13}C this is liquid tetramethylsilane.

To compute the chemical shifts we just need to calculate the current induced by the external magnetic field

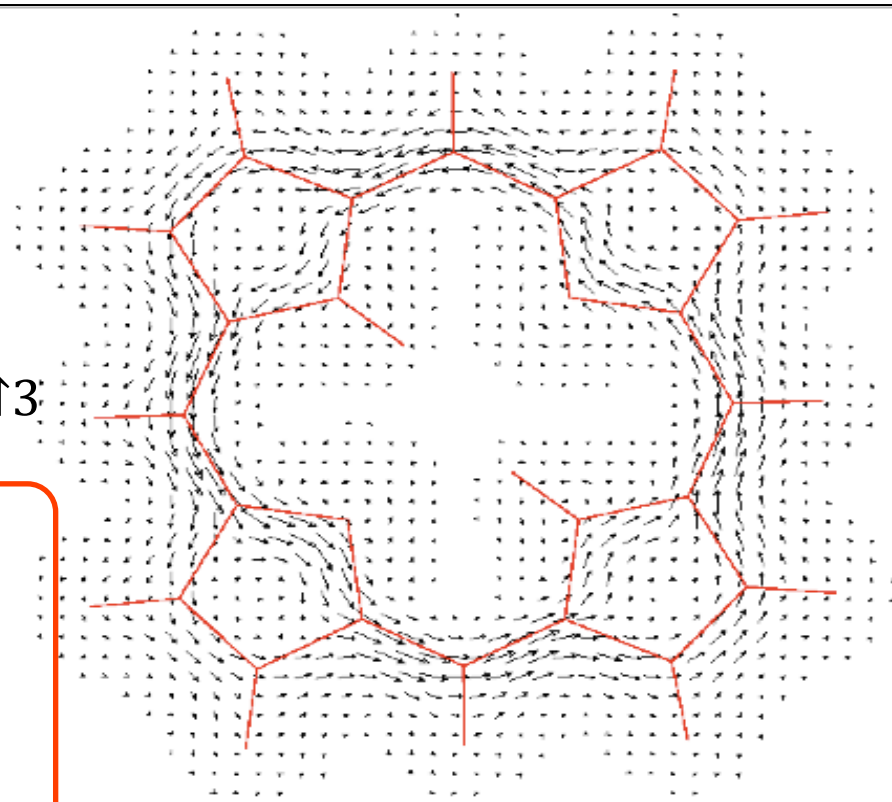
Biot-Savart: **induced magnetic field is**

$$\mathbf{B}_{\text{in}}(\mathbf{r}) = \frac{1}{c} \int d^3\mathbf{r}' \mathbf{j}(\mathbf{r}') \times \frac{\mathbf{r} - \mathbf{r}'}{|\mathbf{r} - \mathbf{r}'|^3}$$

Obtain current within perturbation theory (linear response)

$$O = O^{(0)} + O^{(1)} + O(\mathbf{B}^2)$$

spin contribution averages to zero
in a diamagnetic insulator



$$\mathbf{B} = (1 - \sigma) \mathbf{B}_0$$

$$\mathbf{B}_{\text{in}} = -\sigma \mathbf{B}_0$$

note: σ is a rank 2 tensor

Crystal Structure
X-ray, Neutron diffraction
Cambridge Structural Database

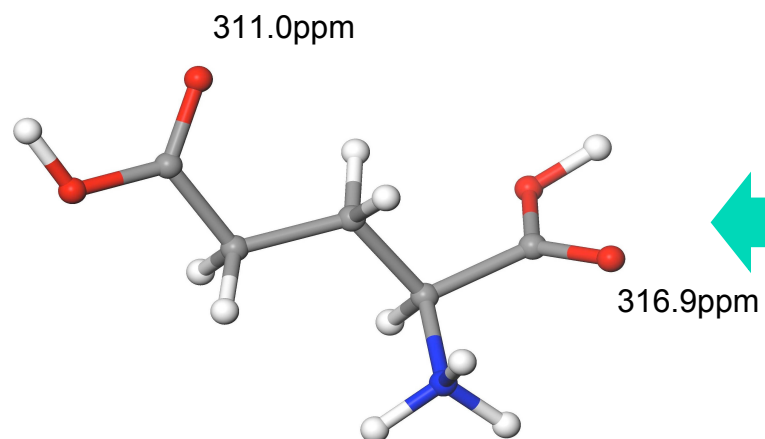
C1	-0.602	-0.537	6.384
C2	0.135	0.622	5.689
H1	0.257	0.361	4.639
...			



Calculate forces
refine structure

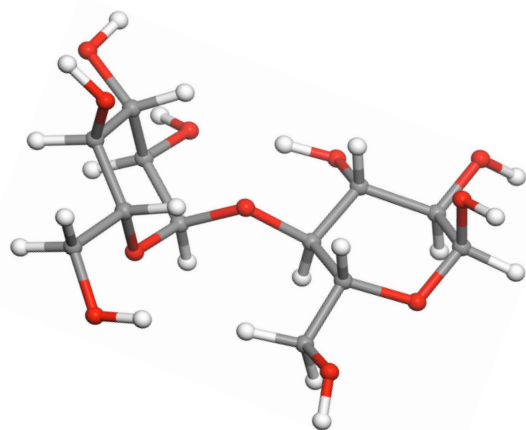


Calculate chemical
shifts

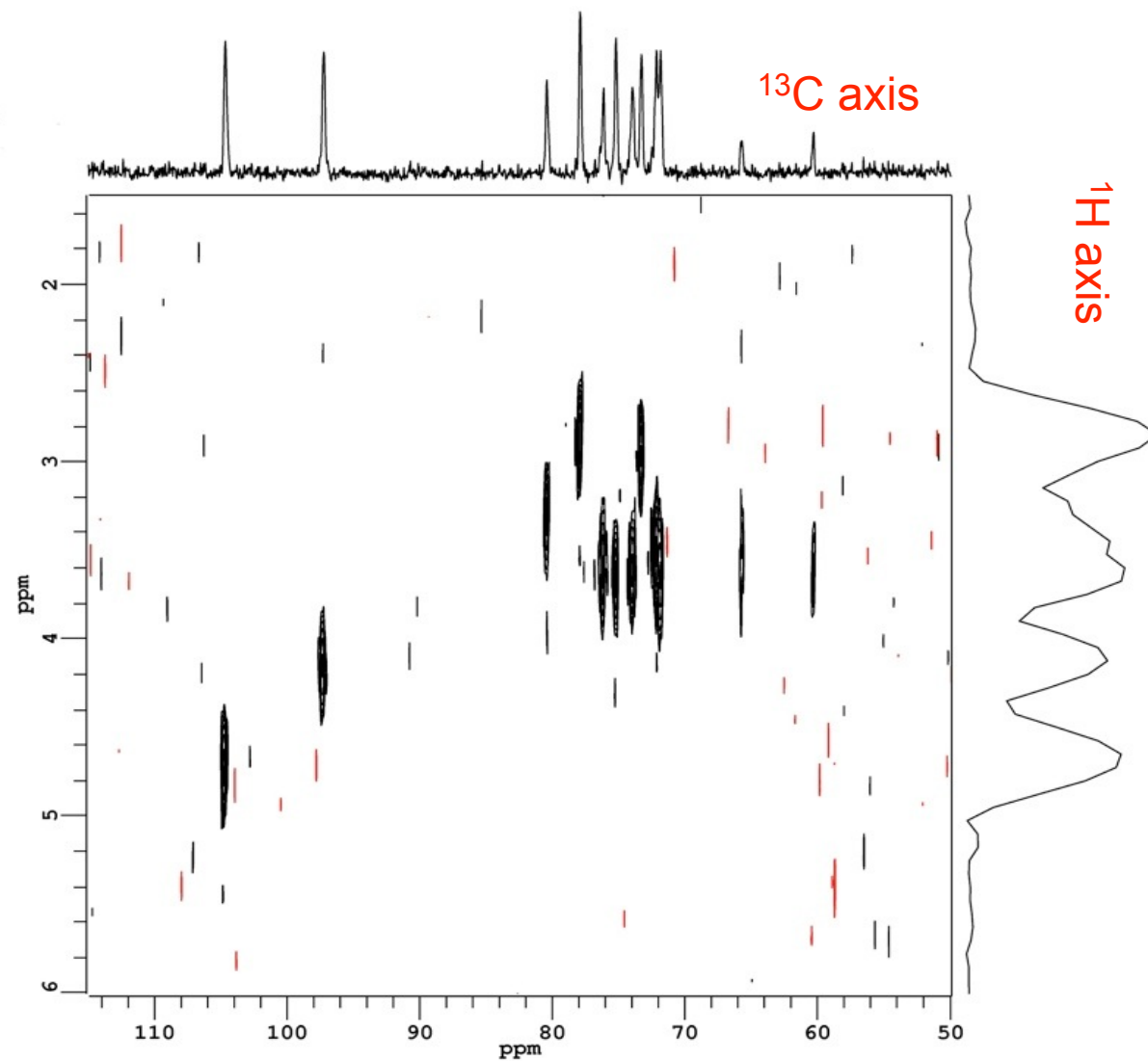


Maltose

sugar used in brewing

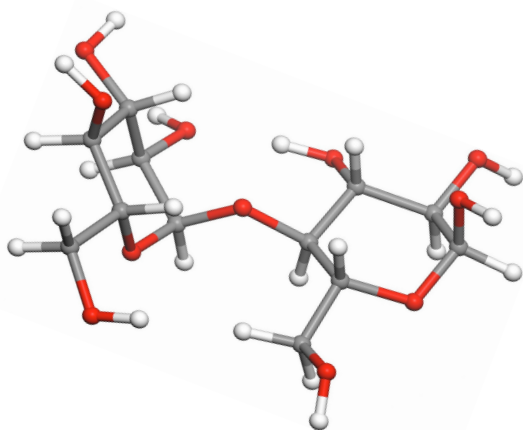


MAS-J-HMQC

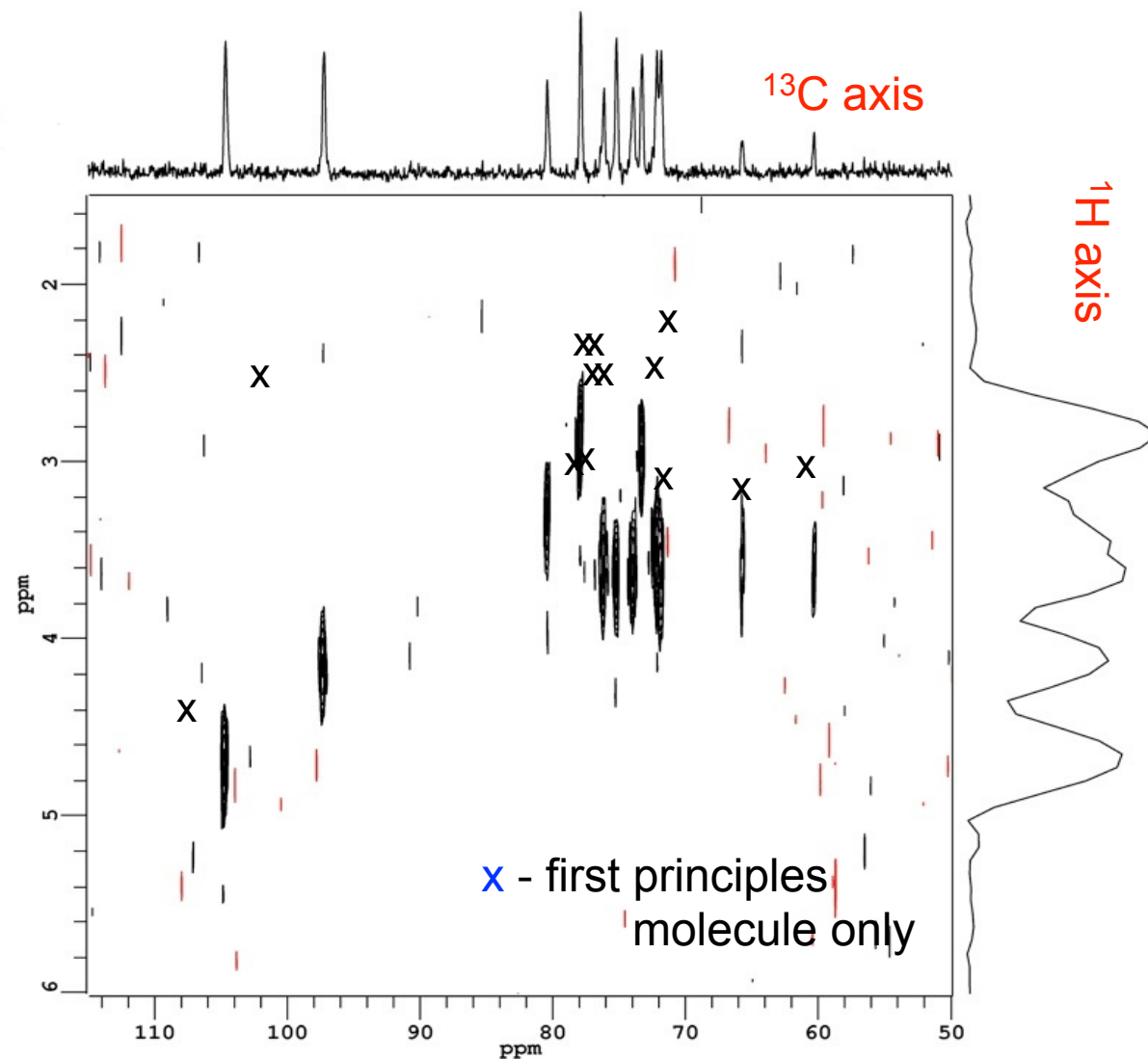


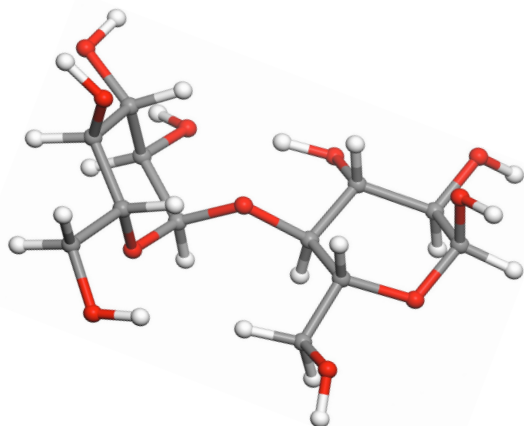
Maltose

sugar used in brewing



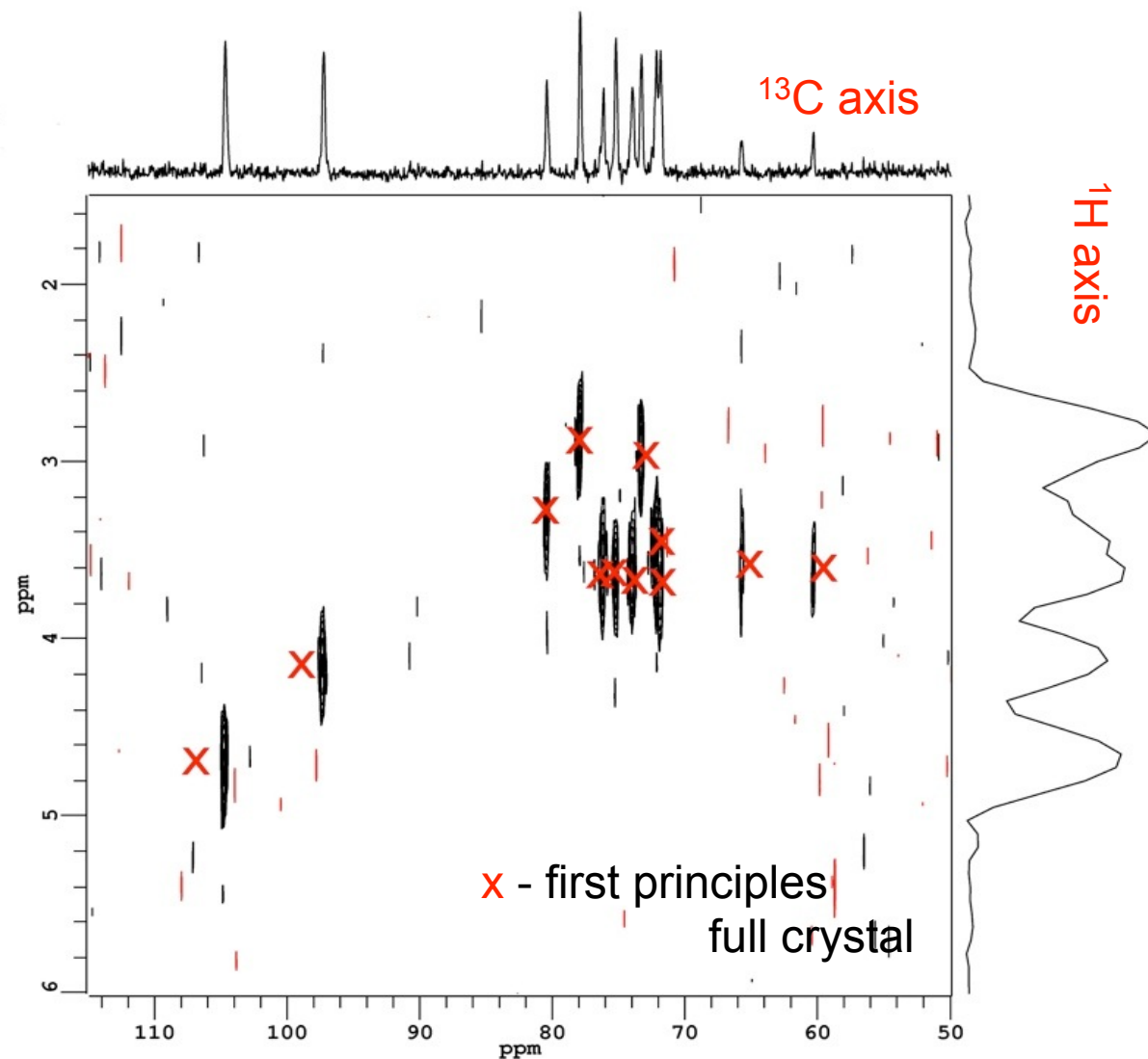
MAS-J-HMQC





Molecule to solid
variation due to
intermolecular
interactions (weak
hydrogen bonds)

J. Am. Chem. Soc. 127 10216
(2005)






Can create a magnetic field in two ways:

1. Intrinsic spin of electrons (i.e. a magnetisation density)
2. Charge of electrons (moving charge gives magnetic field)

Induced Field

$$\mathbf{B}_{\text{in}}^{(1)}(\mathbf{R}_K) = \frac{\mu_0}{4\pi} \int \mathbf{m}^{(1)}(\mathbf{r}) \cdot \left[\frac{3\mathbf{r}_K \mathbf{r}_K - |\mathbf{r}_K|^2}{|\mathbf{r}_K|^5} \right] d^3\mathbf{r} + \frac{\mu_0}{4\pi} \frac{8\pi}{3} \int \mathbf{m}^{(1)}(\mathbf{r}) \delta(\mathbf{r}_K) d^3\mathbf{r}$$

Spin Dipolar  + $\frac{\mu_0}{4\pi} \int \mathbf{j}^{(1)}(\mathbf{r}) \times \frac{\mathbf{r}_K}{|\mathbf{r}_K|^3} d^3\mathbf{r}$  Orbital

 Fermi Contact

$$H_{eff} = - \sum_K \gamma_K \mathbf{I}_K (1 + \sigma_K) \mathbf{B}_{ext} + \sum_{K < L} \mathbf{I}_K (\mathbf{D}_{KL} + \mathbf{J}_{KL}) \mathbf{I}_L$$

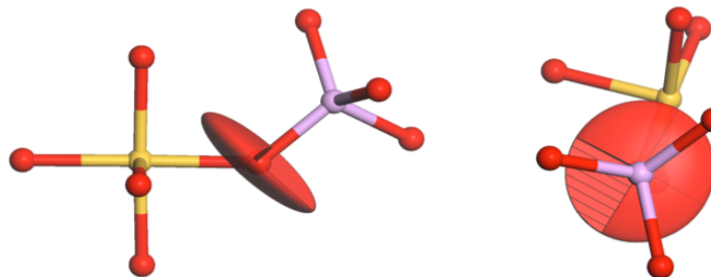
NMR experiments aim to control nuclear spins (I) and external field (B) to measure NMR interaction tensors

We aim to calculate NMR tensors via the quantum mechanics of electrons

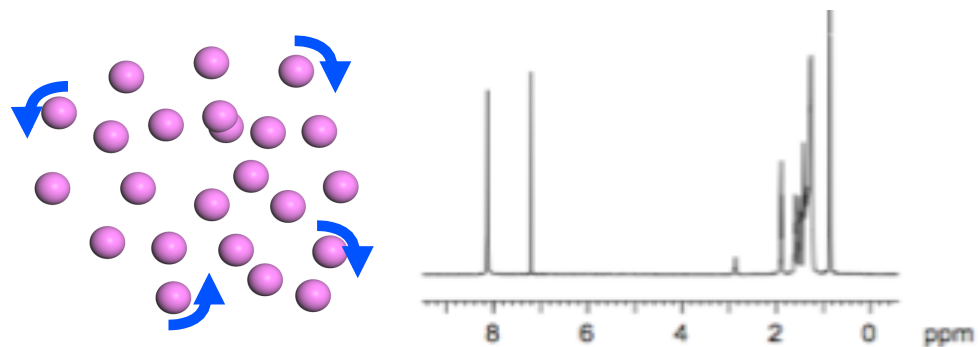
Spectral simulation via nuclear spin dynamics codes

All are tensor properties

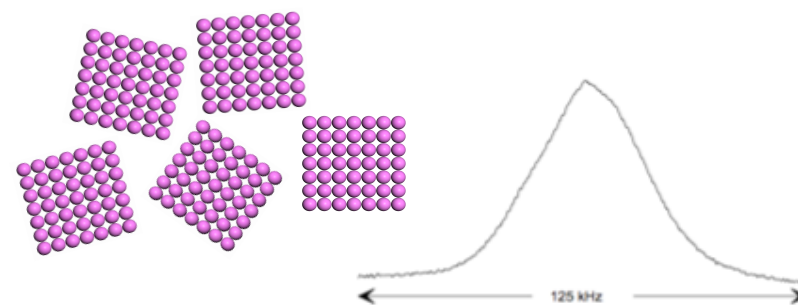
$^1J_{PO}$



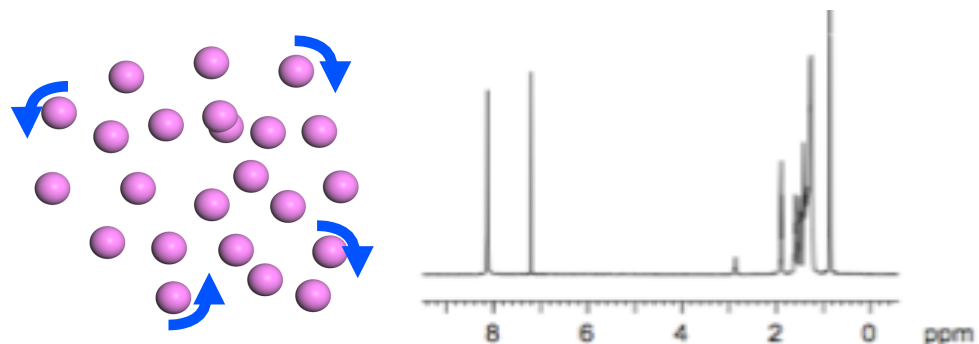
Liquid



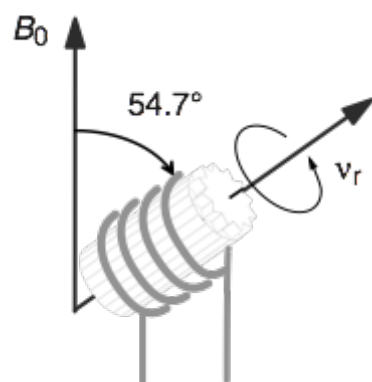
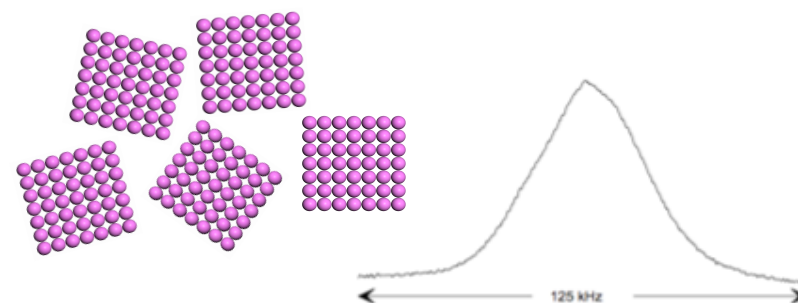
Powdered Solid



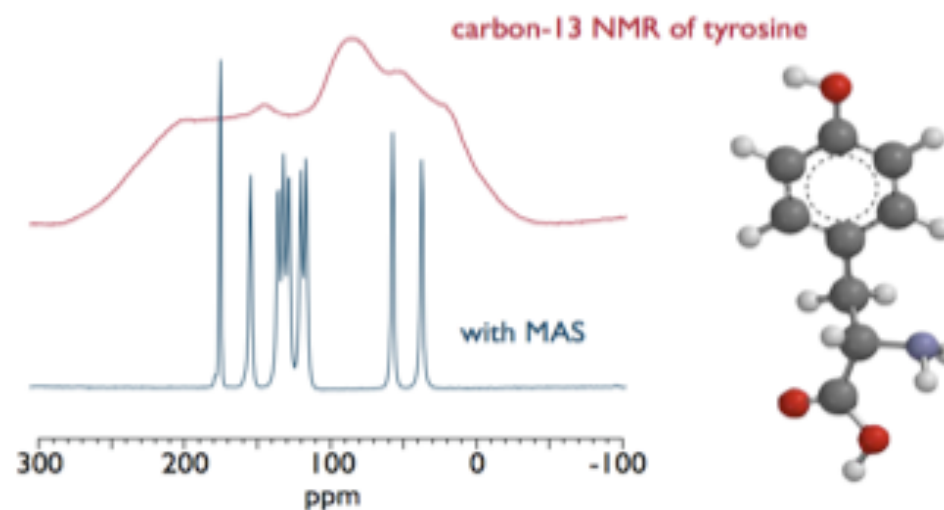
Liquid



Powdered Solid



Anisotropic interactions have first order term proportional to $(3\cos^2\theta-1)$



Chemical Shift

orbital currents create magnetic field

Direct Dipolar Coupling

magnetic field created by neighbouring nuclei
n.b. function of atom positions only i.e. not directly an electronic property

Quadrupolar coupling

nuclei with $I > 1/2$ interact with gradient of electric field (non-spherical charge density)

spin-spin coupling (J-coupling)

magnetic field induced by neighbouring nuclei but mediated via valence electrons

Simulation of Observed Spectra

In general not a simple function of NMR interaction tensors. Need to consider experimental conditions (which NMR experiment). Several sophisticated codes to handle this e.g. SIMPSON, Spinevolution (virtual spectrometers)

Effective Hamiltonian

$$H = -\frac{1}{2\pi} \sum_K \gamma_K I_K (1 - \sigma_K) \mathbf{B}_{\text{ext}}$$

Bare
Dressed

Magnetic Shielding

$$\sigma_K = \frac{\partial^2 E}{\partial \mu_K \partial \mathbf{B}_{\text{ext}}}$$

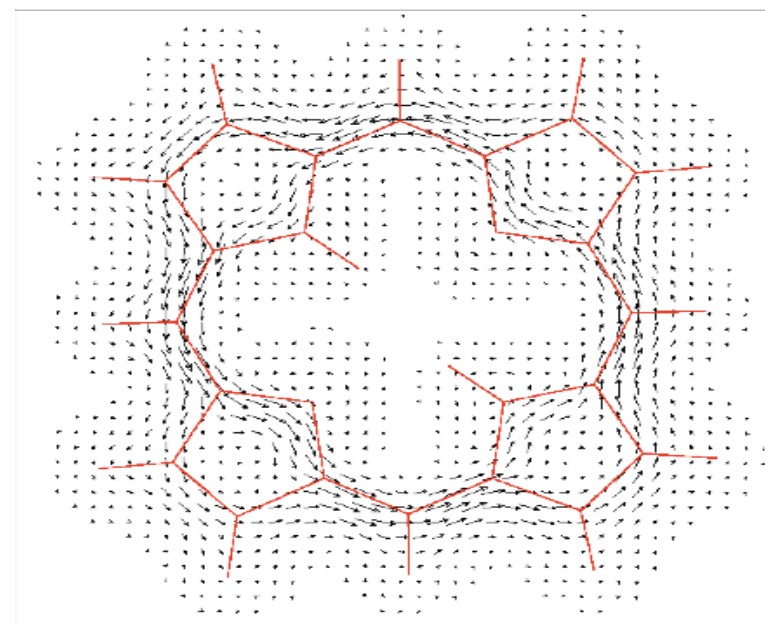
Electronic Response

Induced Current $\mathbf{j}(\mathbf{r})$ Spin Polarisation $\mathbf{m}(\mathbf{r})$
both cause induced magnetic field at nucleus

Induced Field

$$\mathbf{B}_{\text{in}}(\mathbf{R}_K) = \frac{\mu_0}{4\pi} \int \mathbf{j}^{(1)}(\mathbf{r}) \times \frac{\mathbf{r}_K}{|\mathbf{r}_K|^3} d^3\mathbf{r}.$$

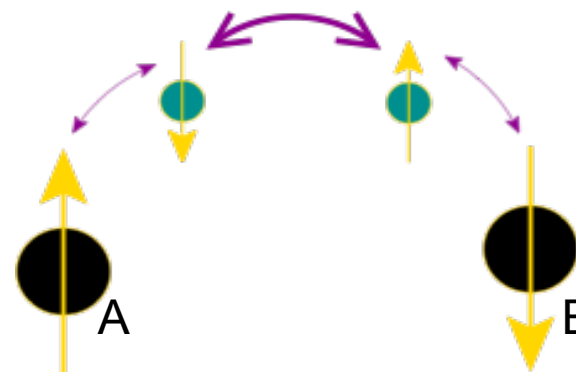
$$\sigma_K = -\frac{\mathbf{B}_{\text{in}}(\mathbf{R}_K)}{\mathbf{B}_{\text{ext}}}$$



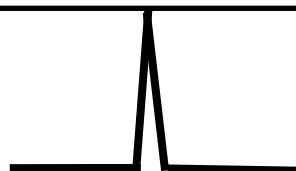
spin contribution averages to zero
in a diamagnetic insulator

Electron-mediated interaction of nuclear spins

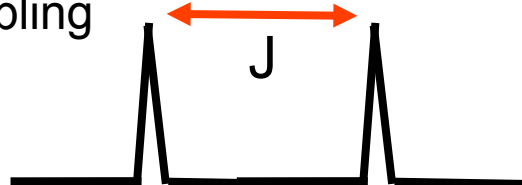
Nucleus A causes a magnetic field at nucleus B (and vice versa)



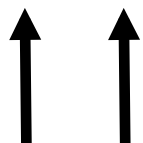
No J-coupling



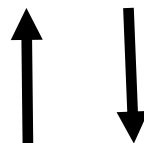
with J-coupling



parallel



anti-parallel



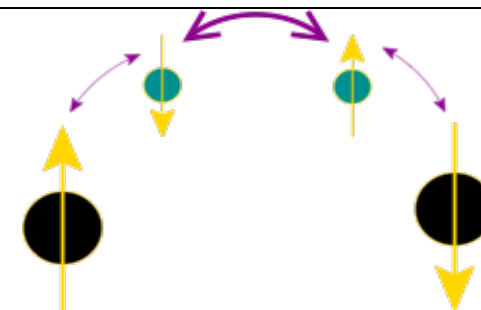
In solids J is rarely revealed in splitting of peaks in 1D spectra (anisotropic interactions broaden peaks)

Increasing use of techniques, eg those using spin-echo modulation, to measure J in solids.

Effective Hamiltonian

$$H = \sum_{K < L} \mathbf{I}_K (\mathbf{D}_{KL} + \mathbf{J}_{KL}) \mathbf{I}_L$$

Bare
Dressed



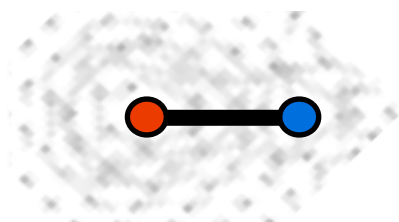
Indirect Coupling

$$\mathbf{J}_{KL} = \frac{\hbar \gamma_K \gamma_L}{2\pi} \frac{\partial^2 E}{\partial \boldsymbol{\mu}_K \partial \boldsymbol{\mu}_L}$$

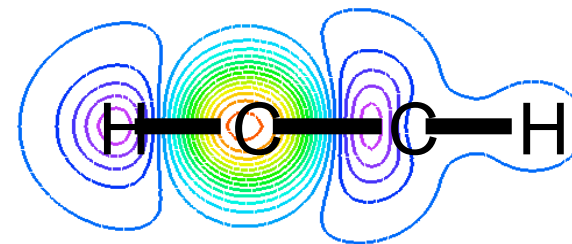
(note E here omits the nuclear-nuclear magnetic interaction)

$$\mathbf{B}_{\text{in}}^{(1)}(\mathbf{R}_K) = \frac{2\pi}{\hbar \gamma_K \gamma_L} \mathbf{J}_{KL} \cdot \boldsymbol{\mu}_L$$

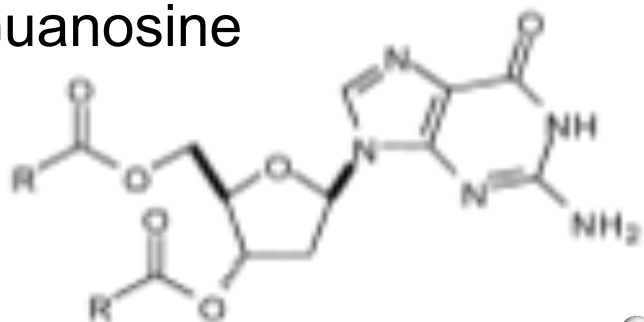
Orbital Magnetic dipole induces orbital motion of electrons

**Orbital Current**

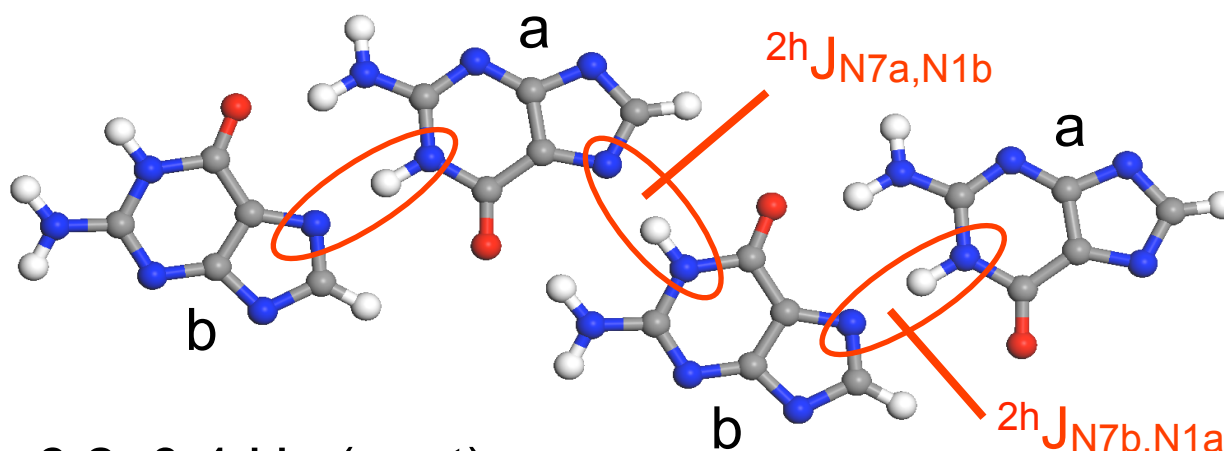
Spin Magnetic dipole induces spin density

**Induced Spin density**

Guanosine



self-assembles into ribbons
molecular electronics (FET)




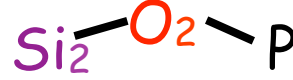


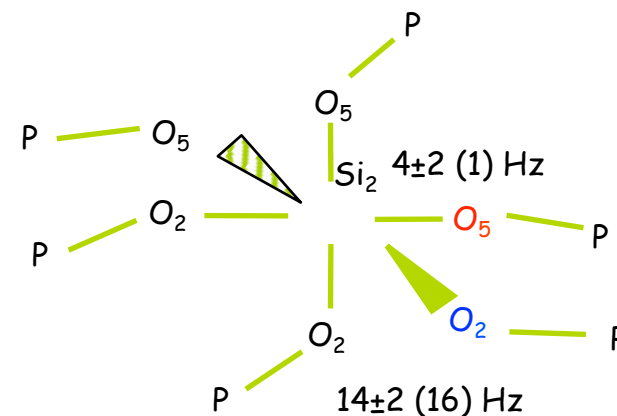
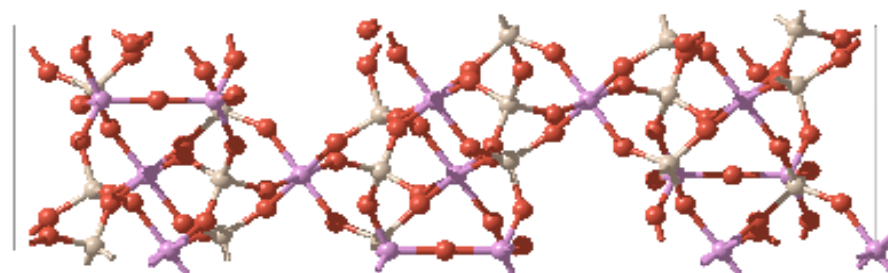
$${}^2hJ_{N7b,N1a} = 6.2 \pm 0.4 \text{ Hz (expt)}$$
$$6.5 \text{ Hz (calc)}$$

$${}^2hJ_{N7a,N1b} = 7.4 \pm 0.4 \text{ Hz (expt)}$$
$$7.7 \text{ Hz (calc)}$$

J. Am. Chem. Soc. **130**, 12663 (2008)



$^2J_{\text{SiP}}$	Expt	Theory
	15 ± 2 Hz	17 Hz
	14 ± 2 Hz	16 Hz
	12 ± 2 Hz	14 Hz
	4 ± 2 Hz	1 Hz



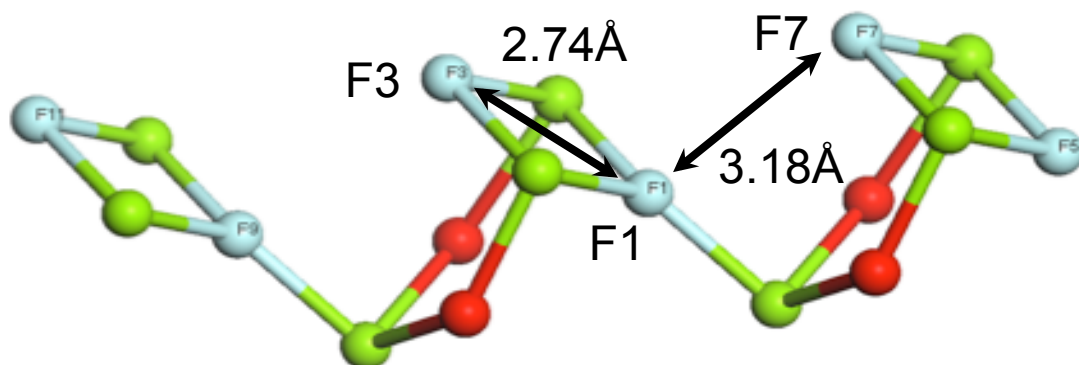
Expt: *Inorg. Chem.* 46, 1379 (2007)

Calc: *J. Chem. Phys.* 127, 204107 (2007)

Design and optimisation of NMR experiments
Interpretation of experiment

Through-space J-coupling

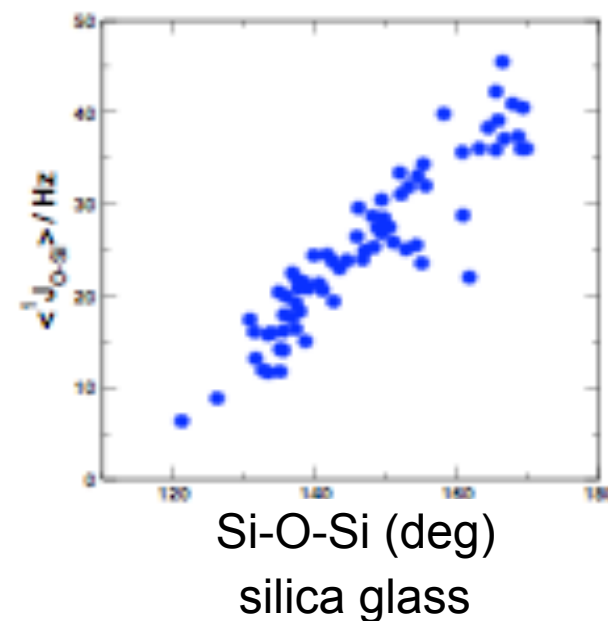
$^1J_{FF}$	Expt	Calc
F1-F3:	18 Hz	12.5 Hz
F1-F7:	2.6 Hz	3.6 Hz



fluorine-substituted deuterated hydrous
magnesium silicate

J. Am. Chem. Soc. 132 15651 (2010)

Distributions of J in amorphous solids



Function of the charge density - ie ground-state property. Also computed by all-electron codes such as Wien2k, Crystal

$$V_{\alpha\beta}(\mathbf{r}) = \int d\mathbf{r}' \frac{n(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|^3} [\delta_{\alpha\beta} - 3(r_{\alpha} - r'_{\alpha})(r_{\beta} - r'_{\beta})/|\mathbf{r} - \mathbf{r}'|^2]$$

charge density

Eigenvalues

$$V_{xx}, V_{yy}, V_{zz} \quad |V_{zz}| > |V_{yy}| > |V_{xx}|$$

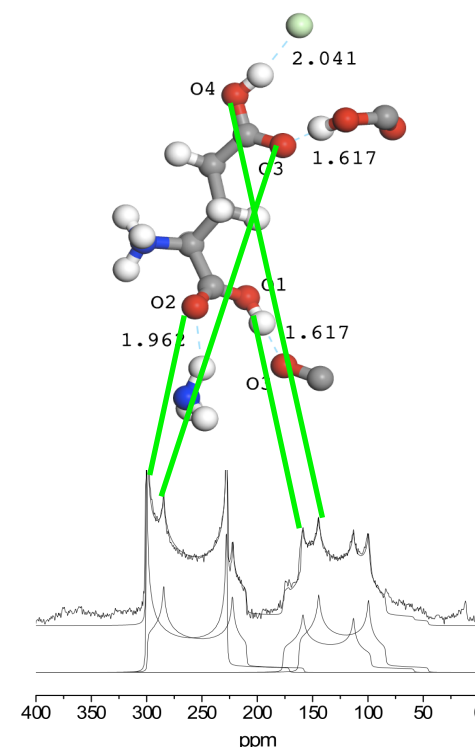
Quadrupolar Coupling

$$C_Q = \frac{eQV_{zz}}{h}$$

Asymmetry

$$\eta_Q = \frac{V_{xx} - V_{yy}}{V_{zz}}$$

Note: The quadrupolar moment, Q, is a nuclear property. Most recent values given in “*Year-2008 Standard Values of Nuclear Quadrupole Moments*”: P. Pyykkö, Mol. Phys. 106, 1965-1974 (2008)
But note Q appears as a simple scaling factor



¹⁷O MAS Glutamic Acid . HCl

*.param file

```
task      : magres  
magres_task : shielding  
          efg  
          nmr  
          jcoupling
```

chemical shift/shielding

electric field gradient

both

Must use on-the-fly pseudopotentials

Highly sensitive to geometry – particularly important to optimise H positions; many structures are obtained from XRD, and H is nearly invisible to X-rays!

CONVERGE!

(basis cut-off & k-points)

Chemical Shielding Tensor					
Nucleus		Shielding tensor			
Species	Ion	Iso(ppm)	Aniso(ppm)	Asym	
H	1	23.81	5.27	0.40	
H	2	24.75	-3.35	0.85	
H	3	27.30	-5.79	0.90	
O	5	-43.73	504.95	0.47	
O	6	-63.53	620.75	0.53	
O	7	-43.73	504.95	0.47	
O	8	-63.53	620.75	0.53	

Anisotropy

Asymmetry

$$\sigma_{\text{aniso}} = \sigma_{\text{zz}} - 1/2(\sigma_{\text{xx}} - \sigma_{\text{yy}}) \quad \eta = 3(\sigma_{\text{yy}} - \sigma_{\text{xx}})/2\sigma_{\text{aniso}}$$

=====

Atom: O 1

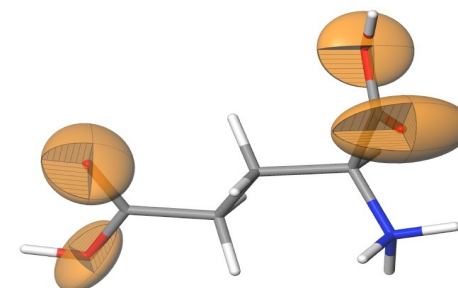
=====

O 1 Coordinates 1.641 1.522 5.785 A

TOTAL Shielding Tensor

218.1858	12.1357	-25.7690
13.4699	191.6972	-7.2419
-25.9178	-6.5205	216.3180

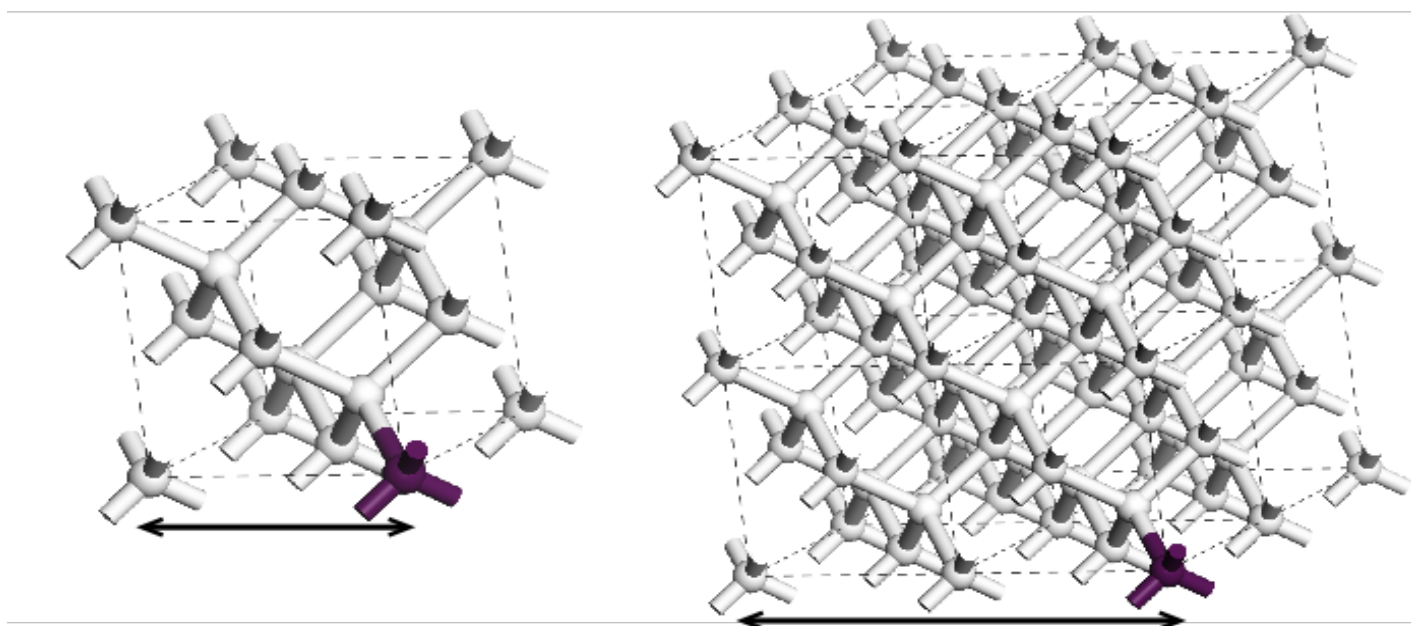
O	1	Eigenvalue	sigma_xx	185.6127	(ppm)	
O	1	Eigenvector	sigma_xx	0.5250	-0.8103	0.2603
O	1	Eigenvalue	sigma_yy	193.8979	(ppm)	
O	1	Eigenvector	sigma_yy	0.4702	0.5310	0.7049
O	1	Eigenvalue	sigma_zz	246.6904	(ppm)	
O	1	Eigenvector	sigma_zz	-0.7094	-0.2477	
O	1	Isotropic:		208.7337	(ppm)	
O	1	Anisotropy:		56.9351	(ppm)	
O	1	Asymmetry:		0.2183		



Note: shielding tensor has a symmetric and an antisymmetric component. Typical NMR experiments are only sensitive to the symmetric part. Therefore we only diagonalise the symmetric part of the shielding tensor

A single calculation gives the coupling between one (perturbing) atom and all others. Might need several calculations to get all of the couplings of interest.

Perturbing atom breaks periodicity - if the unit cell is small you might need to build a supercell to inhibit the interaction with periodic images



Some practical things to keep in mind

Structure

NMR is very sensitive to structure!

XRD atom positions are often not sufficient for accurate simulation of NMR parameters.

Dynamics

Atoms move, and spectra are recorded at finite T.

How to account for that in (static) calculations?

Accuracy

How good are our present (approximate) density functionals?

Effects of Relativity

*scalar relativistic effects easy to treat (relativistic_treatment)
spin-orbit interaction is much harder, but may be needed
(e.g. for shifts on light atoms bonded to heavy atoms)*

Localised d/f Electrons

Local Density Approximation (LDA) not enough, need lower self-interaction error.

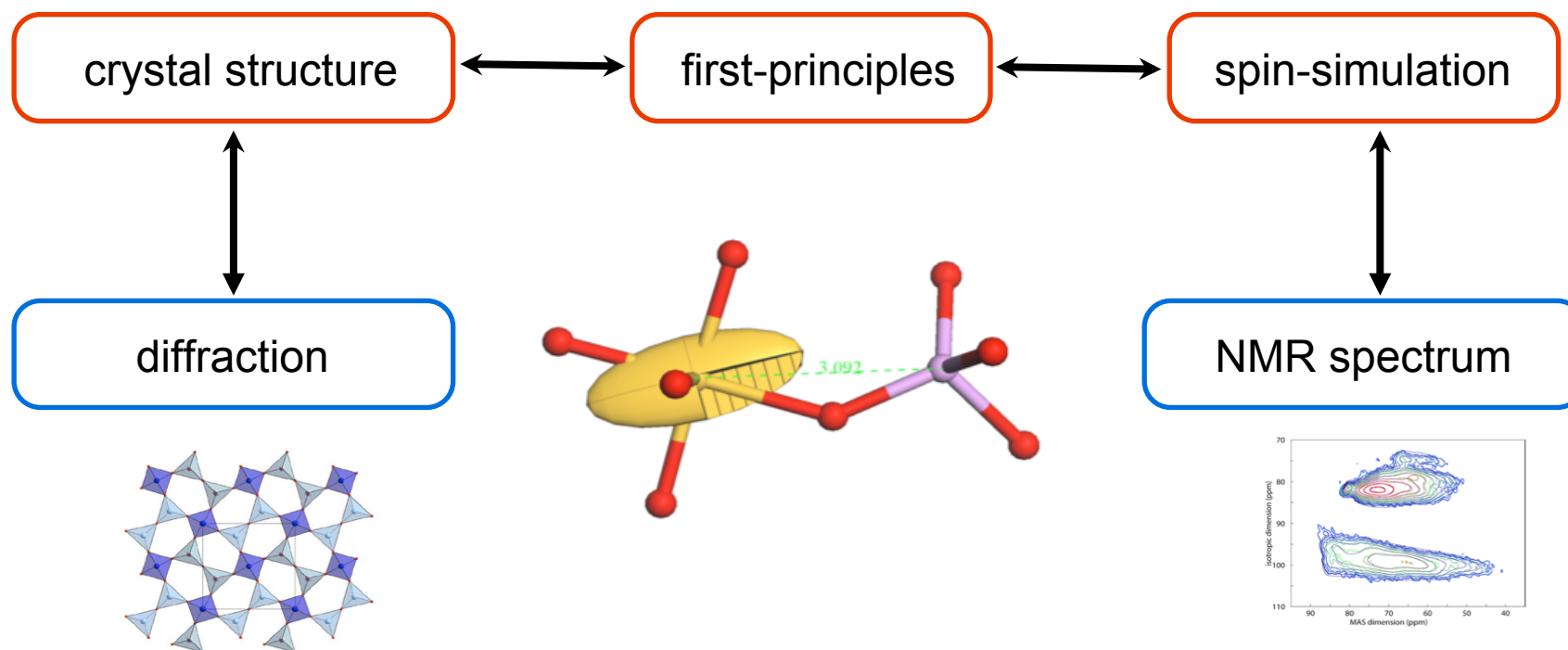
Paramagnetic Materials

Magnetic fields from electronic spin

CCP NC

Collaborative Computational Project for NMR Crystallography

An emerging field, defined as the combined use of experimental solid-state NMR and computation to provide new insight, with atomic resolutions, into structure, disorder and dynamics in the solid state



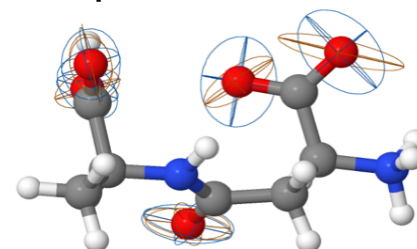
CCP NC

Collaborative Computational Project for NMR Crystallography

Development of Open-source (free) tools for visualisation and processing of first-principles NMR data

Magres Fileformat - community standard format for first-principles NMR data

MagresView - view NMR tensors - link to spin-dynamics

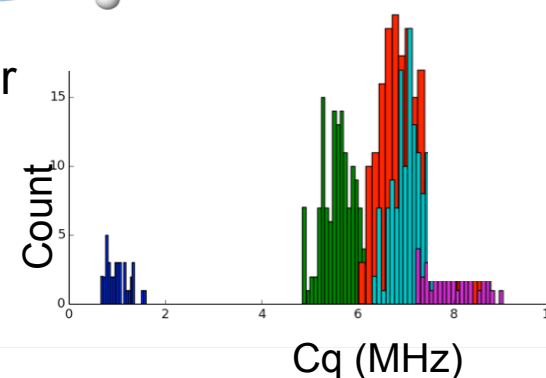


MagresPython - process NMR parameters - dynamics and disorder

www.ccpnc.ac.uk

Validated pseudopotential library - released with CASTEP 8

MagresDatabase - in development



NMR Books

Good Introduction

Nuclear Magnetic Resonance (Oxford Chemistry Primers)

[P. J. Hore](#)

More advanced

Spin Dynamics: Basics of Nuclear Magnetic Resonance

[Malcolm H. Levitt](#)

Solid state NMR

Solid-State NMR: Basic Principles and Practice

David Apperley, Robin Harris, Paul Hodgkinson

Useful survey of applications

Multinuclear Solid-State Nuclear Magnetic Resonance of Inorganic Materials

[Kenneth J.D. MacKenzie](#), [M.E. Smith](#)

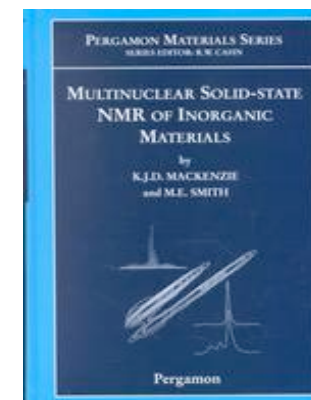
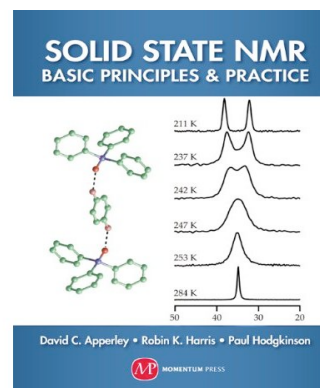
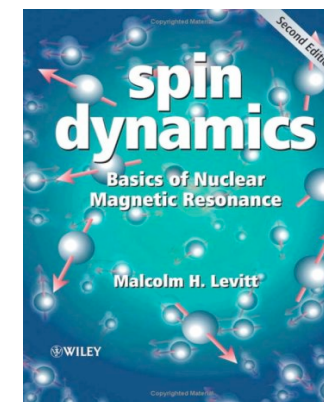
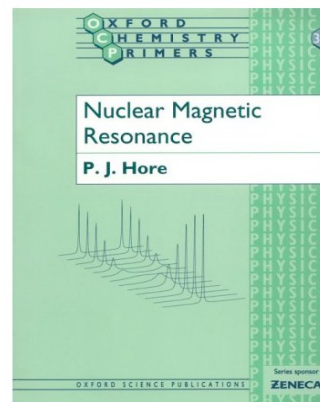
Recent Review Articles

[Recent advances in solid-state NMR spectroscopy of spin \$I = 1/2\$ nuclei](#)

Anne Lesage, *Phys. Chem. Chem. Phys.*, 2009, **11**, 6876

[Recent advances in solid-state NMR spectroscopy of quadrupolar nuclei](#)

Sharon E. Ashbrook, *Phys. Chem. Chem. Phys.*, 2009, **11**, 6892



GIPAW Theory

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