THE UNIVERSITY of York

Structural calculations and phase stability

Matt Probert Condensed Matter Dynamics Group Department of Physics, University of York, U.K. http://www-users.york.ac.uk/~mijp1

- Total Energy
 - Lattice constants
 - Pressure driven phase transitions
 - Bond energy, formation energy, surface ...
- Forces
 - Hellman-Feynman theorem
- Stresses
- Elastic constants



Total Energy

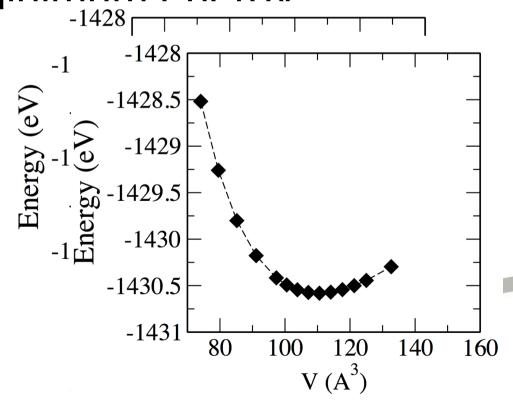


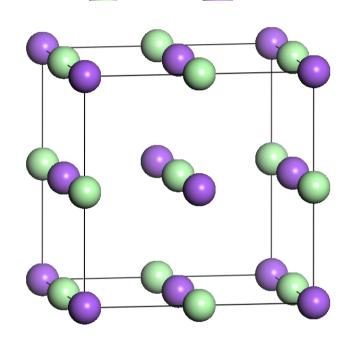
- Quantum Mechanics gives a complete description of materials chemistry & physics
 - QM of electrons and nuclei.
- Covalent, ionic, metallic and van der Waals bonding are all emergent properties of the QM behaviour of electrons.
- Born-Oppenheimer Approximation
 - Treat nuclei classically and assume electrons move adiabatically in a field of fixed ions.
 - QM electrons give interactions between atoms

Total energy

- Primary computed property in DFT is the ground state energy.
 - All observables related to the *change* in energy when the system is perturbed.
- Crystal structure
 - Ion positions and lattice parameters that minimize the total energy.
- Phase stability
 - Phase with lowest E is most stable at T=0
- Elastic constants

- Consider a simple struct.
- The equilibrium CaO latti parameter is given by the minimum energy:







A popular EoS is the Birch-Murnaghan:

$$E(V) = E_0 + \frac{9V_0B_0}{16} \left\{ \left[\left(\frac{V_0}{V}\right)^{\frac{2}{3}} - 1 \right]^3 B'_0 + \left[\left(\frac{V_0}{V}\right)^{\frac{2}{3}} - 1 \right]^2 \left[6 - 4\left(\frac{V_0}{V}\right)^{\frac{2}{3}} \right] \right\}$$

Where the equilibrium volume is V₀ and the bulk modulus is given by

$$B_0 = -V \left. \frac{\partial P}{\partial V} \right|_{P=0}$$

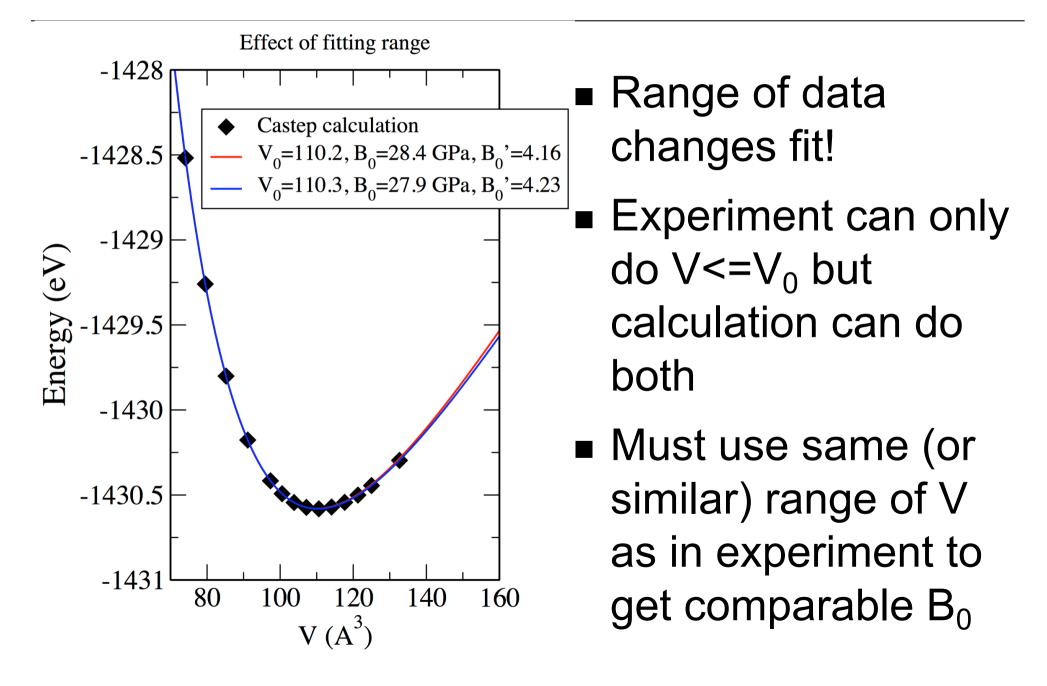
• And the pressure derivative is $B'_0 = \frac{\partial B}{\partial P}\Big|_{P=0}$

Phys. Rev **71** 809 (1947)

- The bulk modulus is a physical (dbservable $B_0' + 16$)
 But two definitions
- - Can extract from fitting B-M EoS to compression data
 - Or from 2nd derivative as in ultrasound J2Dexperiments

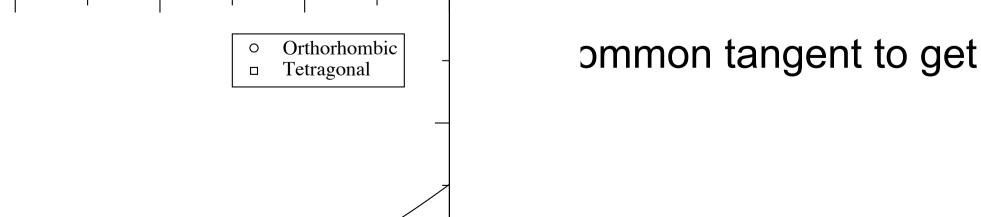
$$B_0 = -V_0 \frac{a}{dV^2}$$

And the results are not the same! Why?



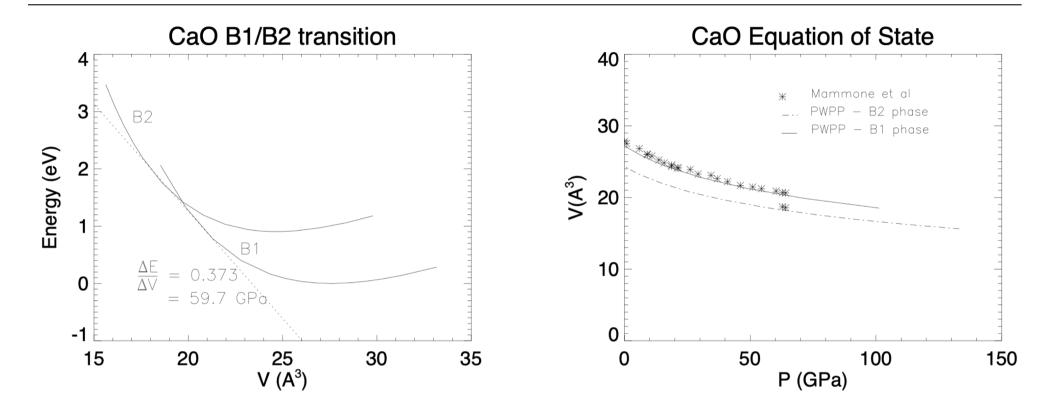
- A structure will change phase to the one of lowest free energy
- Phase equilibrium happens if:
 - Pressure of both phases are equal
 - Enthalpy of both phases are equal:

OTF ultrasofts (PBE)
$$E_2 - E_1 = P(V_2 - V_1)$$



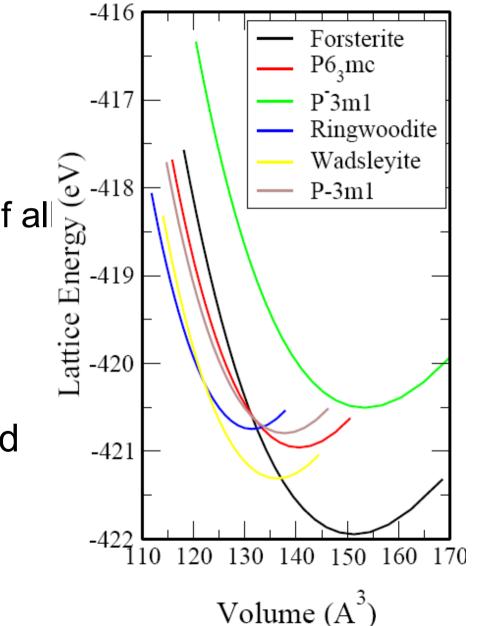


CaO Phase Transition



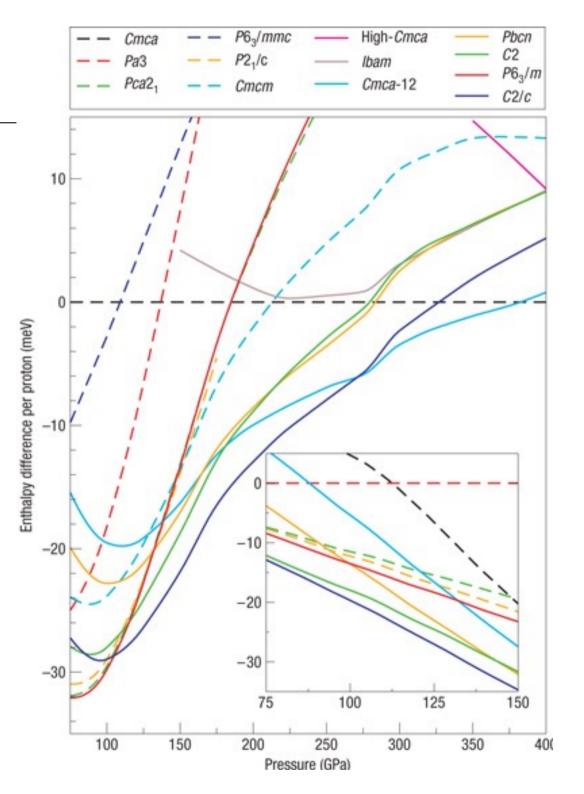
- Or can plot enthalpy H vs pressure P and see which phase has lowest H at given P
- Requires accurate P calculation …

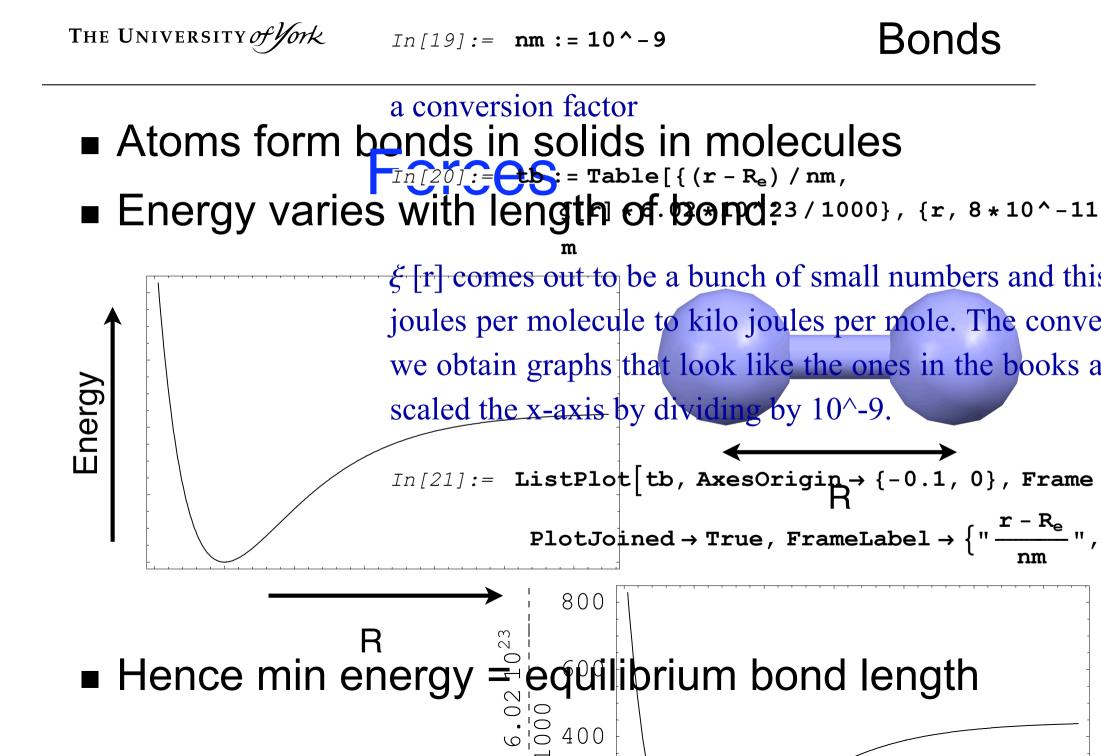
- Which phase is most stable?
 - calculate E(V) curve
 - Forsterite is most stable of al the phases considered
 - then Wadsleyite and then Ringwoodite as increase pressure
- Limited to set of considered structures
 - What if miss something?
 - Active research topic ...



THE UNIVERSITY of York

- High pressure Hydrogen
- Plot ∆H wrt Cmca phase
- Solid lines are new structures, dashed had been already considered by others





- If we stretch a bond too much it will break
 - Energy required is the *dissociation energy*,
 i.e. ∆H = E(molecule) sum E(fragments)
 - Example: methane CH₄
 - dissociation energy = $E(CH_4) E(CH_3) E(H)$
 - Covalent bonds are strong typically 1-5 eV

Environment dependent

- Always requires energy => endothermic
 - H₃C-H bond is ~4.5 eV as bond is short
 - N_2 triple bond is VERY strong = 9.8 eV

- Average dissociation energy over all bonds in a molecule is the *bond energy*
 - More transferable
 - Example: methane CH₄

• bond energy = $\frac{1}{4}$ *(E(CH₄)-E(C)-4E(H))

- NB long bonds are weak
 - BE of H-H = 4.5 eV, length = 0.74 Å
 - BE of H-C = 4.3 eV, length = 1.09 Å
 - BE of H-Si = 3.3 eV, length = 1.48 Å
 - BE of Si-Si = 2.3 eV, length = 2.33 Å

- For solids or liquids
 - The energy required to break all the bonds between the basic units (atoms or molecules) of the material
 - i.e. the intra-molecular bonds
- Example: silicon
 - Cohesive energy = $1/n * (E(Si_n) n*E(Si))$
- Example: water ~ 0.4 eV/molecule
 - Relates to boiling point of water not dissociating water into H₂ and ½ O₂!

- What about defects? Simplest is the point vacancy a single missing atom …
- Neutral vacancy in an elemental solid:

$$E_V = E_{N-1} - \left(\frac{N-1}{N}\right)E_N$$

- Need to consider effect of relaxation around the vacancy – inward or outward?
- Might also be symmetry breaking

More general, for charged defects in nonelemental system:

$$E_{\text{form}} = G_{\text{defect}} - G_{\text{ideal}} = H_{\text{defect}} - H_{\text{ideal}} + \sum_{j} n_{j} \mu_{j}^{\text{removed}} - \sum_{j} n_{j} \mu_{j}^{\text{added}} - n_{\text{e}} \mu_{\text{e}}$$

- Includes change in free energy and change in chemical potential and charge transfer
- Depends on reservoir

THE UNIVERSITY of York

- Where do the defect atoms go/come from?
- Neglects change in entropy

- For an elemental solid, chemical potential μ is the energy of gas phase
- For an compound solid, depends on chemical bonds present

■ E.g. quartz (SiO₂):

$$\mu_{\mathrm{SiO}_2} = \mu_{\mathrm{Si}} + \mu_{\mathrm{O}_2} = \mu_{\mathrm{Si}}^* + \mu_{\mathrm{O}_2}^* + \Delta H_{\mathrm{SiO}_2}$$
$$\mu_{\mathrm{Si}}^* + \Delta H_{\mathrm{SiO}_2} \leqslant \mu_{\mathrm{Si}} \leqslant \mu_{\mathrm{Si}}^*$$

$$\frac{1}{2}\mu_{\mathrm{O}_2}^* + \frac{1}{2}\Delta H_{\mathrm{SiO}_2} \leqslant \mu_{\mathrm{O}} \leqslant \frac{1}{2}\mu_{\mathrm{O}_2}^*$$

where μ_{Si}^* is chem pot of gas-phase Si, and $\mu_{O_2}^*$ is the chem pot of diatomic gas-phase O_2



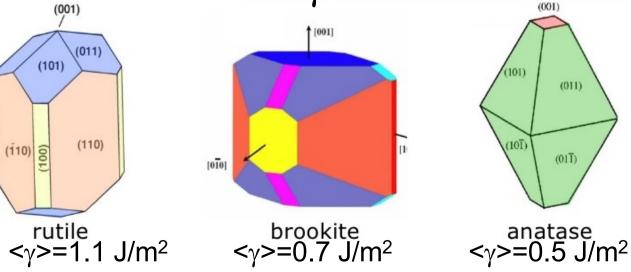
- As well as studying bulk materials and defects, can also study surfaces
- Surface energy γ is the energy required to form a surface: $1 (E = k \cdot E)$

$$\gamma = \frac{1}{2A} \left(E_{surface} - k * E_{bulk} \right)$$

- Where the surface calculation contains 2 surfaces separated by k planes of atoms
- And E_{bulk} is the energy in bulk of 1 plane
- And A is the area of the surface created
- Can also be measured directly in expt

- The equilibrium shape of a crystal or nanoparticle is given by min Gibbs free energy
- Surface Gibbs free energy is $\Delta G = \sum_{i} \gamma_{j} A_{j}$
 - Variational problem minimize △G subject to fixed volume constraint
 - Hence surface with lowest γ will dominate

• e.g. TiO_2

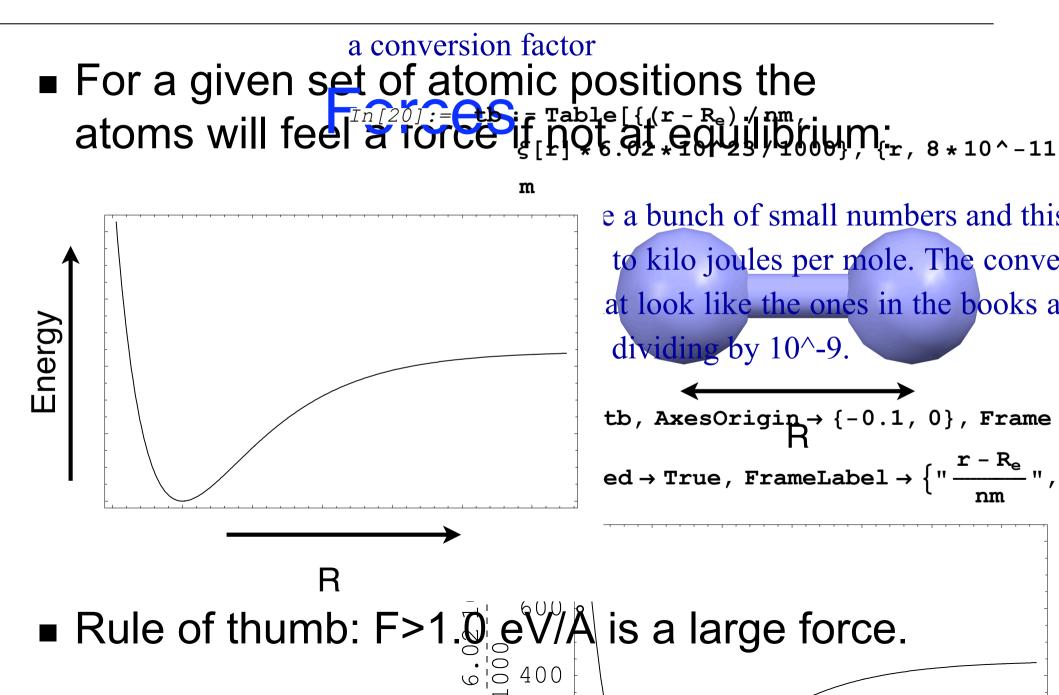




Forces

THE UNIVERSITY of fork $In[19] := nm := 10^{-9}$

Forces



- Minimum energy configuration corresponds to zero force
 - Using forces to find the minimum is much more efficient than just using energy alone
- Minimum enthalpy corresponds to zero force and stress
 - Can therefore minimise enthalpy w.r.t. supercell shape due to internal stress and external pressure
 - Can minimize under external pressure to find new structures

QM Forces

- Classically we have the force **F** at position **R** is determined from the potential energy as $\mathbf{F} = -\nabla_{\mathbf{P}} U(\mathbf{R})$
- Quantum mechanically we therefore expect $\mathbf{F} = -\nabla_{\mathbf{R}} \langle E \rangle$

where

$$\left\langle E \right\rangle = \frac{\left\langle \Psi \left| H \right| \Psi \right\rangle}{\left\langle \Psi \left| \Psi \right\rangle}$$



Stresses

Whilst force is a simple 3-vector, stress is a rank-2 tensor

 $\boldsymbol{\sigma} = \begin{pmatrix} \sigma_{xx} & \sigma_{xy} & \sigma_{xz} \\ \sigma_{yx} & \sigma_{yy} & \sigma_{yx} \\ \sigma_{zx} & \sigma_{zy} & \sigma_{zz} \end{pmatrix}$

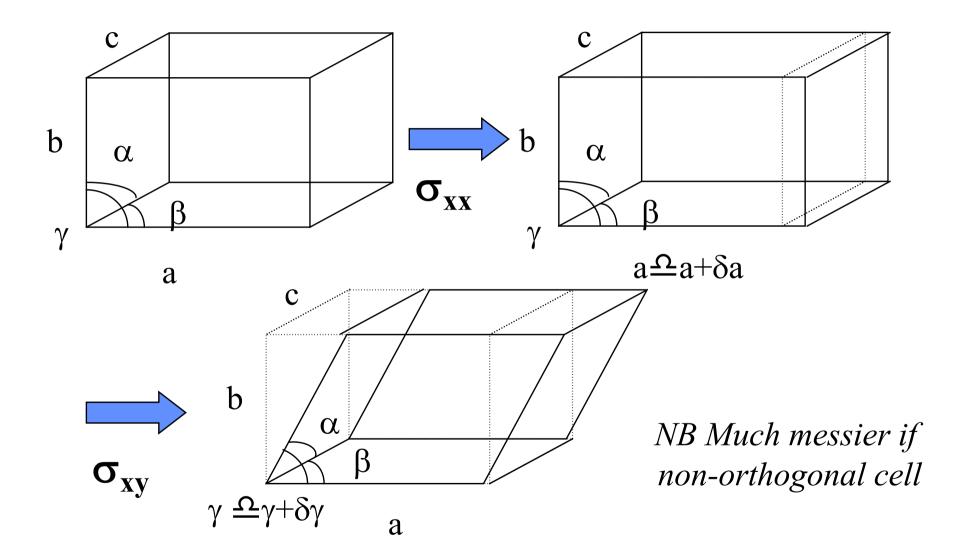
- Symmetry reduces this to 6 independent components:
 - Diagonal = σ_{xx} , σ_{yy} , σ_{zz} = compression

• Off-diagonal = σ_{xy} , σ_{yz} , σ_{zx} = shear

- Convention: P_{ext} + σ = 0 at equilibrium
 - i.e. P>0 = compression, $P_{int}=1/3 Tr(\sigma)$

THE UNIVERSITY of York

Stress in action



If we write the three unit cell vectors a, b, c as the columns of a matrix h then the effect of an applied strain ε is to change the shape of the unit cell:

$$\mathbf{h'} = \left(\mathbf{I} + \boldsymbol{\varepsilon}\right)\mathbf{h}$$

We then have the stress tensor σ related to the strain tensor ε by:

$$\sigma_{\alpha\beta} = \frac{1}{V} \frac{\partial E}{\partial \varepsilon_{\alpha\beta}}$$

where V is the volume of the unit cell.



Elastic Constants

- In 1D you know $E = E_0 + \frac{1}{2}k(x x_0)^2$ where k = elastic constant
- or in terms of dimensionless strain:

$$E = E_0 + \frac{L}{2}k\delta^2$$

In 3D, we have 4th rank elastic constants tensor but can write as 6x6 using Voigt:

$$E(\epsilon) = E_0 + \frac{V}{2} \sum_{ij}^6 C_{ij} \epsilon_i \epsilon_j + O(3)$$

which has max 21 unique components

- For certain high-symmetry crystals there are only a few distinct components to c_{ii}
 - Hence can efficiently calculate with the appropriate minimal pattern of strains
 E.gE (n)a cutsific crystal the cetain of (3) distinct terms: c₁₁, c₁₂ and c₄₄
 - E.g. can choose strain to just study c_{44} :



Summary

Summary

- Total energy is VERY useful
 - Changes can determine phase stability, lattice parameters, bond energies, defect energies, surface energies, etc.
- Can also calculate derivatives
 - E.g. forces (change wrt ionic positions)
 - E.g. stress (change wrt cell parameters)
 - E.g. elastic constants (2nd order strain derivative)

- RM Martin, "Electronic Structure: basic theory and practical methods", Cambridge University Press (2004)
- SJ Clark, MD Segall, CJ Pickard, PJ Hasnip, MIJ Probert, K Refson and MC Payne, *"First principles methods using CASTEP"*, Zeitschrift für Kristallographie **220**, 567 (2005)
- MIJ Probert and MC Payne, "Improving the convergence of defect calculations ...", Phys Rev B 67, 075204 (2003)
- MIJ Probert, "An ab initio study of xenon retention in αquartz", J.Phys Cond Matt 22, 025501 (2010)
- PJ Hasnip, K Refson, MIJ Probert, JR Yates, SJ Clark and CJ Pickard, "Density Functional Theory in the solid state", Phil. Trans. R. Soc. A 372, 20130270 (2014)